

ALMA and *Herschel* observations of the prototype dusty and polluted white dwarf G29-38

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ABSTRACT

ALMA Cycle 0 and *Herschel*¹ PACS observations are reported for the prototype, nearest, and brightest example of a dusty and polluted white dwarf, G29-38. These long-wavelength programmes attempted to detect an outlying, parent population of bodies at 1–100 au, from which originates the disrupted planetesimal debris that is observed within 0.01 au and which exhibits $L_{\text{IR}}/L_* = 0.039$. No associated emission sources were detected in any of the data down to $L_{\text{IR}}/L_* \sim 10^{-4}$, generally ruling out cold dust masses greater than 10^{24} – 10^{25} g for reasonable grain sizes and properties in orbital regions corresponding to evolved versions of both asteroid and Kuiper belt analogues. Overall, these null detections are consistent with models of long-term collisional evolution in planetesimal discs, and the source regions for the disrupted parent bodies at stars like G29-38 may only be salient in exceptional circumstances, such as a recent instability. A larger sample of polluted white dwarfs, targeted with the full ALMA array, has the potential to unambiguously identify the parent source(s) of their planetary debris.

Key words: stars: abundances – circumstellar matter – stars: individual: (G29-38) – planetary systems – white dwarfs.

1 INTRODUCTION

1.1 Discovery and characterization of G29-38

More than a quarter century has passed since the discovery of infrared excess emission from the nearby white dwarf G29-38 (Zuckerman & Becklin 1987). Photometric observations conducted at the NASA Infrared Telescope Facility atop Mauna Kea in the *K*, *L*, and *M* bands revealed flux in excess of that expected for the relatively cool white dwarf. While a warm dust disc was considered a possibility, the infrared excess was initially attributed to a spatially unresolved brown dwarf. In particular, 1000 K circumstellar dust was considered unlikely due to rapid dissipation from radia-

tion drag. Prophetically, Zuckerman & Becklin (1987) noted that if material were orbiting sufficiently close to achieve such a high temperature, then spectral signatures of accretion should be seen.

Within a few years, and thanks to intense observational and theoretical interest from a diverse set of researchers, evidence began to disfavour a substellar companion as the origin of the infrared emission. First, some of the earliest infrared imaging arrays revealed G29-38 to be a point source in all available bandpasses. Secondly, near-infrared spectroscopy measured a thermal continuum (Tokunaga et al. 1988), whereas a very cool atmosphere was expected to exhibit absorption features. Thirdly, the detection of optical stellar pulsations echoed in the near-infrared was difficult to reconcile with a brown dwarf secondary (Graham et al. 1990; Patterson et al. 1991). Fourthly and finally, significant 10 μm emission was detected at G29-38, at a level a few times greater than that expected for an object with a Jupiter-sized radius, essentially ruling out the brown dwarf companion hypothesis (Telesco, Joy & Sisk 1990; Tokunaga, Becklin & Zuckerman 1990).

A decade after the discovery of its infrared excess, the optical and ultraviolet spectroscopic detection of multiple metal species in

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the atmosphere of G29-38 (Koester, Provencal & Shipman 1997) made it clear that the star was accreting from its immediate environs. Although white dwarfs with metal absorption features have been known for nearly a century (vMa 2; van Maanen 1917), the source of the heavy elements had never been observationally identified and the closely orbiting disc at G29-38 was the smoking gun.

1.2 Tidally destroyed planetesimals

In a seminal paper, Jura (2003) modelled the observed properties of G29-38 by invoking a tidally destroyed minor planet (i.e. large asteroid) that evolves into an opaque, flat ring of dust analogous to the rings of Saturn. The particles are heated by the star, producing an infrared excess, and slowly dragged down on to the stellar surface, polluting its otherwise pristine atmosphere with heavy elements. The tidally disrupted asteroid model has seen continued success since its inception, and is considered the standard model for metal-enriched white dwarfs. In the intervening decade, an enormous amount of observational progress has occurred (for a detailed review, see Farihi 2011), all of which supports the accretion of asteroid-like debris in dynamically active, post-main-sequence planetary systems (Veras et al. 2013).

(1) Over 30 metal-lined white dwarfs are now known to exhibit $T \sim 1000$ K thermal emission from discs (e.g. Kilic et al. 2006; Jura, Farihi & Zuckerman 2007; von Hippel et al. 2007; Farihi, Jura & Zuckerman 2009; Farihi et al. 2010b, 2012; Brinkworth et al. 2012; Girven et al. 2012; Xu & Jura 2012); their properties are precisely as expected for material contained within the Roche limit of the star and feeding the stellar surface (Bochkarev & Rafikov 2011; Rafikov 2011; Metzger, Rafikov & Bochkarev 2012). A fraction of these exhibit metallic, gaseous emission (Gänsicke et al. 2006, 2008; Gänsicke, Marsh & Southworth 2007; Melis et al. 2011, 2012; Farihi et al. 2012) that is spatially coincident with the particulate discs (Melis et al. 2010; Brinkworth et al. 2012).

(2) All dusty white dwarfs observed with IRS on *Spitzer* exhibit strong silicate emission features consistent with olivines (Reach et al. 2005, 2009; Lisse et al. 2008; Jura, Farihi & Zuckerman 2009a), and which are also seen in the infrared spectra of evolved solids associated with planet formation (Lisse et al. 2008).

(3) The elemental abundances in disc-polluted white dwarfs are universally depleted in volatile elements (especially carbon and hydrogen) and have a refractory rich pattern that broadly mimics the terrestrial material of the inner Solar system (Zuckerman et al. 2007; Klein et al. 2010; Gänsicke et al. 2012; Jura & Young 2014). All data acquired to date are consistent with rocky parent bodies that formed interior to a snow line (Jura & Xu 2012; Farihi, Gänsicke & Koester 2013), including evidence for differentiation (Farihi et al. 2011; Zuckerman et al. 2011; Gänsicke et al. 2012; Jura, Xu & Young 2013).

(4) The collective properties of all published ($N \sim 250$ as of this writing), metal-polluted white dwarfs support the accretion of planetary debris for the population at large (Zuckerman et al. 2003, 2010; Farihi et al. 2010a), including even those stars that do not exhibit detectable infrared excesses (Rocchetto et al. 2014). These findings imply that at least 30 per cent (Koester, Gänsicke & Farihi 2014) of cool white dwarfs have the signatures of large planetesimals, whose size estimated diameters range from ~ 10 to ~ 1000 km (Jura & Young 2014; Wyatt et al. 2014).

1.3 Motivation for long-wavelength data

All successful models for discs at white dwarfs invoke tidally destroyed planetary bodies that originate in a more distant and substantially more massive reservoir of planetesimals (Debes & Sigurdsson 2002; Jura 2003; Bonsor, Mustill & Wyatt 2011; Debes, Walsh & Stark 2012; Frewen & Hansen 2014). To date, there have been no indications of cooler dust associated with an outlying planetesimal population around metal-polluted white dwarfs, but the longest wavelength observations conducted for a substantial number of stars are both *WISE* 22 μm and *Spitzer* MIPS 24 μm photometry (Jura et al. 2007; Farihi et al. 2009; Hoard et al. 2013). These data only probe for dust at radii of 1 to a few au for typical 10 000–20 000 K white dwarfs, and thus longer wavelength data are needed to search for populations analogous to the asteroid and Kuiper belts of the Solar system.

This paper presents *Herschel* (Pilbratt et al. 2010) far-infrared and Atacama Large Millimeter/Submillimeter Array (ALMA) observations of G29-38 (ZZ Psc, WD 2326+049). This iconic stellar remnant has a hydrogen-dominated atmosphere of $T_{\text{eff}} \approx 12\,000$ K, intrinsic brightness $\log(L/L_\star) \approx -2.5$, and a cooling age near 380 Myr (Fontaine, Brassard & Bergeron 2001; Giammichele, Bergeron & Dufour 2012). At 13.6 pc, G29-38 is the nearest example of a disc-polluted white dwarf and also the brightest by an order of magnitude, making it ideal for long-wavelength observations sensitive to cold dust emission, and where a ring of 20–200 au diameter would span 1.4–14 arcsec on the sky and be potentially resolved.

Both sets of observations resulted in null detections and provide limits on cold dust masses similar to known Kuiper belt objects. The observations and data analysis details are presented in Section 2, from which are derived sensitivities to fractional dust luminosity, and similarly to dust masses, as a function of temperature and orbital radius for all existing $\lambda \geq 70$ μm observations. The results are presented in Section 3 together with a comparison of these data and sensitivities with known dusty A-type stars, which represent possible progenitors of the G29-38 system.

2 OBSERVATIONS AND DATA

2.1 *Herschel*

G29-38 was targeted by the *Herschel Space Observatory* on 2012 June 6 with the Photodetector Array Camera and Spectrometer (PACS; Poglitsch et al. 2010) at 100 and 160 μm . These data are particularly sensitive to 20 K dust in the 10 au region, where any emission at these wavelengths is expected to be unresolved at the 7 arcsec/100 μm diffraction limit of a 3.5 m telescope. While circumstellar dust orbiting beyond 50–80 au can be spatially resolved around nearby stars with *Herschel* (e.g. Booth et al. 2013), the very low luminosity of a cool white dwarf like G29-38 makes such a detection unlikely (see Section 3.1).

The source was expected to be substantially fainter than the 50 mJy limit recommended for standard chop-nod observations, and thus the mini-scan map mode was used. These were executed in three pairs of cross-scans, using individual cross-scan angles of 70° and 110°, and each of the six segments repeated 25 times for an on-source time of 1800 s. In total, G29-38 was observed for 3.0 h on source and with 9.4 h observatory time.

The PACS data were reduced using the *Herschel* Interactive Processing Environment (HIPE) version 7.0. The data were processed

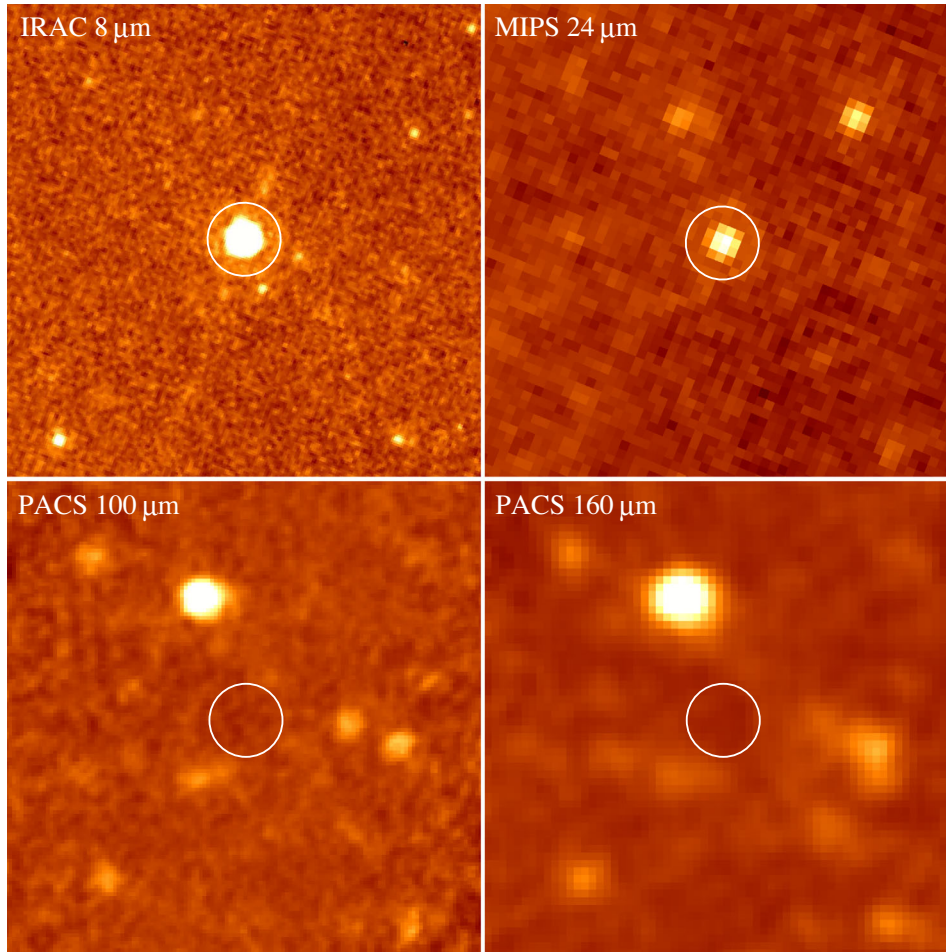


Figure 1. Reduced infrared images of G29-38 taken with both *Spitzer* and *Herschel*. The images are north up and east left and approximately 125 arcsec on a side. Accounting for its proper motion of $0.49 \text{ arcsec yr}^{-1}$, the expected location of G29-38 is marked by a circle in each image, with emission only detected in the *Spitzer* data.

using the standard PACS photometer processing steps, and maps were made using the `PHOTPROJECT` task. The data were high-pass filtered in the scan direction with filter width of 66 and 102 arcsec (equivalent to 16 and 25 frames) in the 100 and 160 μm bands, respectively. The fully reduced images are shown in Fig. 1 alongside *Spitzer* imaging detections in the mid-infrared.

The source was assumed to be point-like and photometry was performed using apertures. Source apertures with radii of 5 and 8 arcsec ($0.7 \times \text{FWHM}$, which is near optimal for a point source) were used for the 100 and 160 μm bands, respectively. The field background and 1σ noise were obtained by taking the median and standard deviation of multiple apertures, with radii equal to that of the source apertures, located within 1 arcmin of the source location, respectively. The measured fluxes ($F_{100,160\mu\text{m}} = -0.5, -0.6 \text{ mJy}$), and sum in noise apertures ($\sigma_{100,160\mu\text{m}} = 0.7, 1.6 \text{ mJy}$), have all had aperture corrections applied; 1.94 and 1.90 for the 100 and 160 μm bands, respectively. The derived upper limits for G29-38 are the flux measured in the source apertures $+3\sigma$, and are listed in Table 1.

2.2 ALMA

G29-38 was also observed with the ALMA as part of Early Science operations at the beginning of Cycle 0 in 2011 November.

Data were acquired over the course of several nights, with two 17 min observations on 18 and 26 November in band 6 (230 GHz, 1305 μm), and five 23 min observations on 6, 14, and 16 November in band 7 (345 GHz, 870 μm). During the brief time span of the observations, the proper motion of G29-38 was negligible for the concatenation of data sets. All but two observations were made using 15 antennas, but a single observation in each of bands 6 and 7 was made with 17 and 14 antennas, respectively. The spatial configuration of the array provided baselines covering a range from 10 to 150 m.

Initial calibrations and pointing were done on the bright quasars 3C 454.3 and 3C 446, while Neptune was used for absolute flux calibration. The science observations were interleaved with the phase calibrator J2323–032. The observing setup was the most sensitive for continuum observations, using the wide-band TDM mode with 128 spectra channels of 15 625-MHz wide. 2 MHz bandwidth per polarization, and resulting in a full effective bandwidth of 7.5 MHz.

The data were processed by the ALMA pipeline and the achieved rms values were $0.071 \text{ mJy beam}^{-1}$ in band 6 and $0.056 \text{ mJy beam}^{-1}$ in band 7; these were adopted as the 1σ noise values. No sources were detected in either band, and upper limits for unresolved emission from the science target were taken to be 3σ and are listed in Table 1. Using the Common Astronomy Software Applications

Table 1. Multiwavelength fluxes and upper limits for G29-38.

Source	λ_{eff} (μm)	F_{ν} (mJy)
GALEX	0.15	3.4
GALEX	0.23	10.9
U	0.37	17.3
B	0.44	21.7
V	0.55	22.3
R	0.64	18.8
I	0.80	15.1
J	1.24	8.9
H	1.66	6.0
K_s	2.16	5.6
IRAC	3.55	8.4
IRAC	4.49	8.8
IRAC	5.73	8.4
IRAC	7.87	8.4
IRTF	10.5	11.1
IRS	16.0	3.7
MIPS	23.7	2.4
MIPS	71.4	<1.3
PACS	100	<1.6
PACS	160	<4.4
ALMA	870	<0.17
ALMA	1305	<0.21

Notes. Fluxes are *GALEX* far- and near-ultraviolet (Martin et al. 2005), ground-based *UBVRIJHK_sN* (Tokunaga et al. 1990; Skrutskie et al. 2006; Landolt & Uomoto 2007), *Spitzer* IRAC, IRS 16 and MIPS 24 μm photometry (Reach et al. 2005; Farihi, Zuckerman & Becklin 2008). Also listed are 3σ upper limits from MIPS 70 μm photometry (Jura et al. 2009b), and the similar PACS and ALMA limits reported here.

package, a band 6 and 7, beam-combined image was created to increase the chance of source detection using the increased signal and spectral information. No emission from the science target is evident in these merged down to $3\sigma \approx 0.17$ mJy. The beam-combined image is displayed in Fig. 2 and is overlaid with the band 7 beam size (1.5 arcsec \times 1.8 arcsec at position angle 1:63).

3 ANALYSIS AND DISCUSSION

3.1 Sensitivity to dust emission

Fig. 3 details the spectral energy distribution of G29-38, extending from the ultraviolet to submillimetre wavelengths. The stellar photosphere is detected from the far-ultraviolet until the near-infrared, where the spatially unresolved emission of the inner disc dominates at $\lambda > 2 \mu\text{m}$. Of the suite of available data, the ALMA observations had the best ability to spatially resolve any emission from dust, but only for relatively cold and distant material beyond 11 au.

Following Wyatt (2008), all the available infrared and submillimetre upper limit F_{ν} at G29-38 were transformed into limiting fractional disc luminosities as a function of dust temperature and corresponding orbital radius using

$$f = 3.4 \times 10^9 F_{\nu} d^2 X_{\lambda} / r^2 B_{\nu}(\lambda, T), \quad (1)$$

where d is the distance to the star, r the orbital radius of dust grains of temperature T , and B_{ν} is the Planck function. The factor X_{λ} allows for the faster falloff in emission towards longer wavelengths for non-blackbody – smaller, warmer yet more distant – grains (Williams & Andrews 2006). These upper limit fractional dust luminosities are

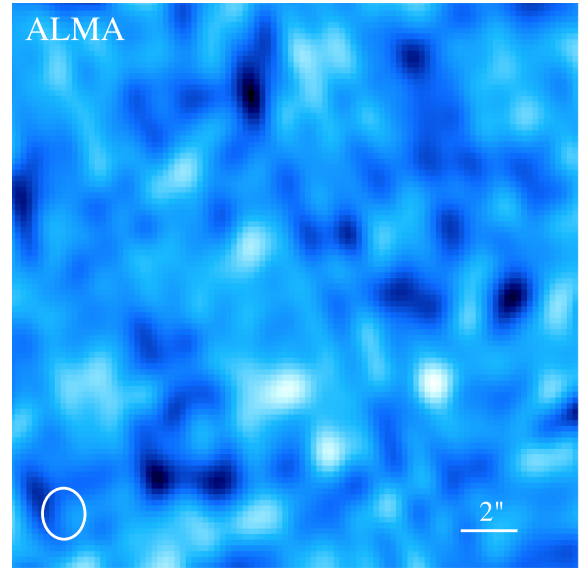


Figure 2. Bands 6 and 7, beam-combined ALMA image centred on the expected position of G29-38. The image is north up and east left, with the band 7 beam size and orientation shown in the lower left, and the spatial scale given in the lower right. No sources are detected above $3\sigma \approx 0.17$ mJy.

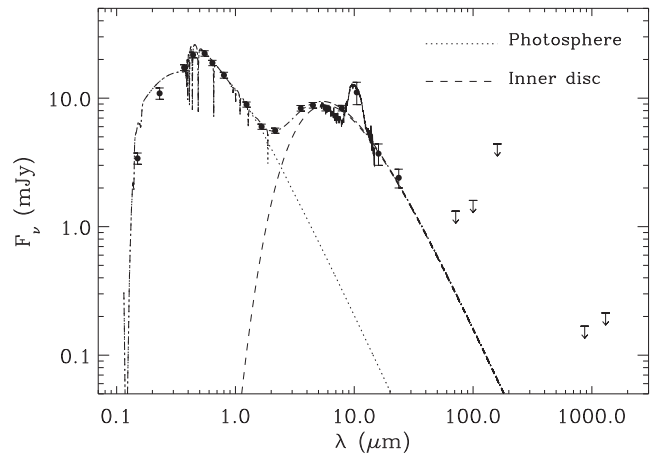


Figure 3. Extended spectral energy distribution of G29-38. Details of the measured fluxes and 3σ upper limits are listed in Table 1 and discussed in Section 2. The dotted line is a stellar atmosphere model, the dashed line is a face-on, flat and optically thick ring model ($T_{\text{in}} = 1250$ K, $T_{\text{out}} = 650$ K) fitted to the warm disc continuum emission, and the solid line the measured strong silicate feature in the *Spitzer* IRS SL spectrum (Reach et al. 2009).

plotted as coloured curves in Fig. 4 for the *Spitzer*, *Herschel*, and ALMA observations. Also included in the plot is a similar curve for the most compact configuration of the Expanded Very Large Array (EVLA) at 45 GHz, with a 1.5 arcsec beam and a 5.6 μJy continuum sensitivity (Perley et al. 2011).

Any blackbody dust grains with temperature below approximately 20 K would be orbiting beyond 11 au and hence be spatially resolved to some degree in the ALMA observations, which had a band 7 beam diameter between 1.5 and 1.8 arcsec, and thus capable of resolving ring structures with radii larger than 10.2–12.2 au at 13.6 pc (see Fig. 2). The plotted ALMA sensitivities were corrected for such resolved cases by assuming a face-on disc and dividing by

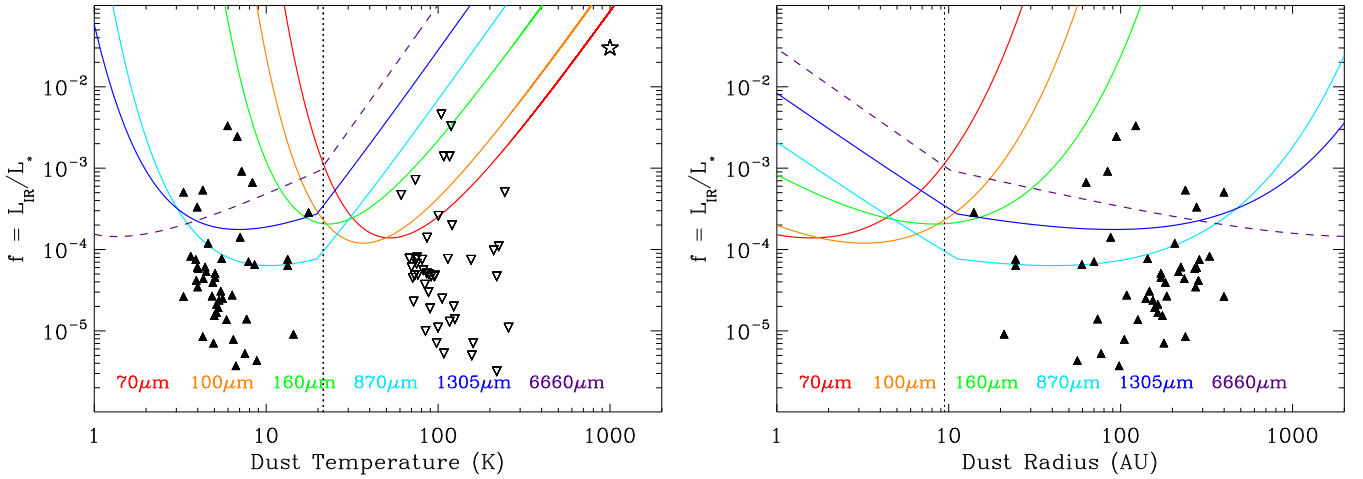


Figure 4. Upper limit sensitivity curves for all long-wavelength non-detections of G29-38 as a function of dust fractional luminosity, temperature, and orbital radius; only the area above each curve was detectable in individual observations. Limits with *Spitzer* MIPS 70 μm are shown in red, *Herschel* PACS 100 and 160 μm in orange and green, respectively, and ALMA 870 and 1305 μm in light and dark blue, respectively. The ALMA sensitivities have been corrected for spatially resolved emission from dust beyond 11.2 au (below roughly 20 K for blackbody grains). Also shown as a dashed purple line are predictions for the EVLA at 45 GHz in the most favourable configuration (Section 3.1). The filled triangles are evolved, cool white dwarf stage projections (Section 3.2) of known, dusty A-type stars (Su et al. 2006), while the progenitor systems are shown as open, inverted triangles in the left-hand panel. The observations as a whole were particularly sensitive to evolved main belt analogues, shown as a dotted line in each plot. The star symbol in the upper-right corner of the left-hand panel is the (undetected) known warm disc, orbiting interior to 0.01 au.

the number of beams per ring circumference. The sensitivities of the *Herschel* observations were similarly corrected for dust located beyond 50–80 au, but it is worth noting that such high fractional luminosities (1) lie above the Fig. 4 plots and cannot be assumed to be optically thin, and (2) would have been detected as spatially resolved emission with ALMA.

3.2 Projections of dusty A-type stars

Also in Fig. 4 are plotted model extrapolations for dusty A-type stars from Su et al. (2006), as their *current* dust properties would appear if the host stars were evolved into cool white dwarfs with parameters similar to G29-38. To begin, the inferred orbital radii of the dust components on the main sequence were expanded by a factor $r_{\text{wd}}/r_{\text{ms}} = M_{\text{ms}}/M_{\text{wd}} \approx 3.5$, appropriate for G29-38 and a representative value for A stars based on the initial-to-final mass relations (e.g. Kalirai et al. 2008; Williams, Bolte & Koester 2009). This change not only leads to a change in dust grain temperatures, but also in their illumination. The amount of light intercepted by the dust decreases by $(r_{\text{ms}}/r_{\text{wd}})^2$, but the total emitting area of the dust *increases* because small grains are not removed by radiation pressure (Farihi et al. 2008). On the main sequence, the collisional cascade is truncated at a grain diameter that is approximately $L_{\text{ms}}/M_{\text{ms}}$ (Artymowicz 1988), and thus the emitting area and fractional luminosity are proportional to $1/\sqrt{L_{\text{ms}}/M_{\text{ms}}}$.

Following Bonsor & Wyatt (2010), considering the realistic emission properties of small grains illuminated by the faint white dwarf, only particles larger than 0.1 μm contribute significantly to the emission, and thus the change in fractional luminosity will be

$$f_{\text{wd}}/f_{\text{ms}} \approx (r_{\text{ms}}/r_{\text{wd}})^2 \times \left(\sqrt{L_{\text{ms}}/M_{\text{ms}}}/\sqrt{0.1} \right). \quad (2)$$

Notably, several of the brightest A star debris disc projections were readily detectable in the ALMA (but not *Herschel* or *Spitzer*) observations, even with the modest number of antennas and corresponding limited sensitivity. However, it is important to note that these

modelled points (1) are based on discs orbiting relatively young, main-sequence stars, (2) essentially preserve the mass of colliding planetesimals between stellar phases, and (3) allow arbitrarily cold dust, including temperatures below that supported by ambient interstellar radiation. A fixed minimum temperature for dust grains that is substantially warmer than the 3 K cosmic microwave background will favourably impact their detectability with ALMA as such a disc will tend towards the more sensitive parts of the curves in Fig. 4, while the assumption of no further collisional evolution or dust depletion (e.g. wind drag or dynamical) over the course of stellar evolution is admittedly physically unrealistic.

While small dust grains would be removed during the asymptotic giant phase, they are quickly replenished by collisions. The collision rate of the largest surviving particles can be long, but the smallest grains are repopulated from the integrated collisions of all large bodies, resulting in a replenishment time-scale equal to their depletion time-scale in a steady state (Wyatt, Clarke & Booth 2011). Still, detailed models for the post-main-sequence evolution of Kuiper belt analogues (Bonsor & Wyatt 2010; Bonsor et al. 2011) suggest that detecting these discs in cool white dwarf systems like G29-38 is challenging, primarily due to collisional depletion of the disc material. Such collisional depletion would occur on even shorter time-scales for asteroid belt analogues favoured by the volatile poor abundance patterns seen via atmospheric pollution (Gänsicke et al. 2012; Jura & Xu 2012), including G29-38 in particular (Xu et al. 2014). And while searches for cold discs have been carried out with *Spitzer* in more favourable, hot (pre-)white dwarf systems with relatively high stellar luminosity (Su et al. 2007; Chu et al. 2011), the resulting picture is complicated by binarity and dusty outflows associated with their immediate progenitors (Clayton et al. 2014).

The stochastic planetesimal accretion models of Wyatt et al. (2014) place G29-38 among the top 1 per cent of accretors in terms of incoming material, but it is also notable among DA stars where instantaneous rates can be inferred from observed metal abundances. If white dwarf pollution were correlated with disc brightness on the main sequence, then G29-38 would have evolved from one

of the brightest main-sequence discs and thus imply the highest chance of cold dust detection among polluted white dwarfs. The fact that it is not detected, and the fact that the young dusty A star discs are predicted to be detectable in the absence of depletion of their belts, may support, but not strongly, the likelihood that such main-sequence discs, on average, become depleted in dust as they progress significantly beyond their current, relatively young ages.

These caveats notwithstanding, the analysis and figures indicate that in the absence of significant depletion, or in the case of where debris is replenished in these orbital regions (e.g. Stone, Metzger & Loeb 2014), ALMA is an excellent tool to both detect and spatially resolve discs at white dwarfs. Giant impacts or instabilities analogous to the late heavy bombardment would substantially increase the detectability of planetesimal discs at white dwarfs, and cold dust can accumulate over long time-scales due to the feeble stellar luminosity. Another possibility that favours the retention of substantial dust mass is if the planetesimals have high eccentricity and thus a reduced collision rate (Wyatt et al. 2010). While the average planetesimal disc at white dwarfs may be depleted, some fraction of the population may remain salient. The final ALMA array of 66 dishes should more than double the sensitivity, all else being equal, and thus be capable of detecting circumstellar dust more representative of current A stars, as well as asteroid belt analogues, now cold and expanded beyond several au.

3.3 Dust mass limits

To derive upper limits on the mass of dust present around G29-38, the 870 μm observations were the most sensitive for realistic temperatures. Emission at this wavelength traces predominantly millimetre-sized dust particles and is most likely optically thin. In the absence of any detailed information on the composition and size of any grains associated with this source, the average value of dust opacity is adopted here, $\kappa = 1.7 \text{ cm}^2 \text{ g}^{-1}$ at 870 μm , noting that the actual opacity may be anywhere in the range 0.2–4.0 $\text{cm}^2 \text{ g}^{-1}$ (Draine 2006).

Using a radial dependence of the dust temperature identical to that derived in Section 3.1 and plotted in Fig. 4, the adopted upper limit flux of $F_\nu = 0.17 \text{ mJy beam}^{-1}$ is converted to the dust mass using

$$m(T) = F_\nu d^2 / \kappa B_\nu(T). \quad (3)$$

A range of radial distances r were explored, where dust emission falls within one half of the primary beam of a 12 m ALMA antenna at 870 μm , or up to 60 au from the star. Having no information on the radial or azimuthal distribution of dust, nor its projected distance in the plane of the sky, the calculations assume a simple, circular, and face-on ring of radial thickness much less than the beam size (22 au at 13.6 pc).

Spatially unresolved emission allows direct conversion from the observed flux to the mass of the dust using the above expression, while at $2r > 22$ au the ring would become resolved and the observed flux per beam is then a fraction of the total flux in the ring. Upper limits derived in this way are shown in Fig. 5, where only high dust masses greater than that of Pluto ($1.3 \times 10^{25} \text{ g}$) can be ruled out beyond 11 au in the spatially resolved regime. In the case of unresolved emission within 11 au, the upper limit dust masses are relatively low and comparable to a few times the mass of Ceres ($9.4 \times 10^{23} \text{ g}$).

Coincidentally, the unresolved dust mass limits fall within a factor of a few of the highest known masses of metals residing in the

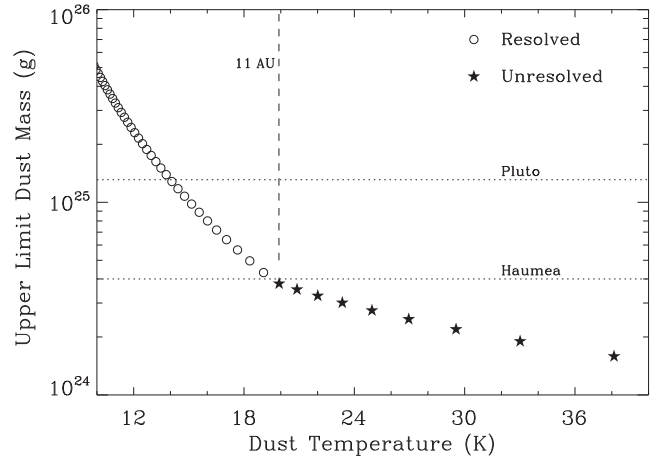


Figure 5. Dust mass limits calculated from the achieved ALMA upper limit at 870 μm of 0.17 mJy. For dust warmer than around 20 K, the emission is expected to remain spatially unresolved in the observations, while cooler dust would have been spread across many beams and thus the sensitivity drops accordingly.

outer layers of polluted white dwarfs with significant convection zones (Dufour et al. 2010; Farihi et al. 2010a). However, there is no reason to expect a physical connection between collisionally generated dust masses in outer planetesimal belts (which requires replenishment over appropriate time-scales) and the amount of mass that is transported via (presumably) intact parent bodies to the tidal disruption radius, and later the stellar surface, of polluted white dwarfs.

4 OUTLOOK

Debris from the inferred parent population of planetesimals at G29-38 remains undetected with *Herschel* PACS at 100 and 160 μm , and ALMA in early and relatively shallow observations at 870 and 1305 μm . While the ALMA observations were best suited to outer regions analogous to evolved Kuiper belts, the *Herschel* data were uniquely sensitive to an evolved, asteroid-like belt at G29-38 and in general to dust at temperatures and orbital regions intermediate to the relatively warm dust seen at polluted white dwarfs and the cooler dust often detected at main-sequence stars. The non-detection at G29-38 is not wholly unexpected, as disc evolution models supported by observations of main-sequence stars predict that the available mass in both dust and parent bodies decreases significantly over time-scales of several hundred Myr (Wyatt et al. 2007), and this depletion is likely enhanced during the post-main sequence (Bonsor & Wyatt 2010).

These data collectively preclude a relatively bright, $L_{\text{IR}}/L_* > 10^{-4}$ disc in the general vicinity of 10 au around the nearest and brightest polluted white dwarf with an infrared excess. From orbital expansion alone, a planetesimal belt currently in within 10 au would have orbited within the terrestrial zone – interior to the water ice line – during the main sequence, and thus be consistent with the suspected source regions for parent bodies of the disc of disrupted and polluting material in G29-38 and a growing number of white dwarfs where detailed abundance measurements allow a robust assessment of their volatile content (Gänsicke et al. 2012; Jura & Xu 2012; Farihi et al. 2013; Koester

et al. 2014; Xu et al. 2014). Therefore, any dusty asteroid-like belt must lie below these detection limits.

However, recent work suggests that it is difficult for sufficient material to survive to the white dwarf phase in this inner region (Debes et al. 2012; Frewen & Hansen 2014). A parent body origin in an outer, Kuiper-like belt remains consistent with these long-wavelength observations (Bonsor et al. 2011), but the volatile deficiency combined with substantial parent body masses, both inferred via atmospheric pollution, remains to be explained.

While significant uncertainty remains, ALMA is the only current facility able to empirically constrain the source regions for the parent bodies of the planetary debris surrounding and falling on to white dwarfs like G29-38.

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