

1 **Structural and electronic characterisation of Cu/Au(111) near-**
2 **surface alloys**

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14 Geometrical structure and electronic characteristics of ultrathin metallic films of one metal
15 on another are strongly influenced by factors such as lattice mismatch and formation of near-
16 surface alloys. Doped systems often show modified chemical-physical properties which may
17 be amenable to different reactivity routes. Copper-gold surface alloys have received the
18 attention of several studies, only a few of which have been undertaken in an UHV
19 environment, using surface sensitive techniques. In this contribution, systems produced upon
20 room temperature deposition of copper on the $(22\times\sqrt{3})$ -Au(111) surface, at various copper
21 loadings and annealed to different temperatures, are investigated using scanning tunnelling
22 microscopy and synchrotron radiation based techniques (X-ray photoelectron diffraction,
23 photoemission), with comparison with some theoretical modelling. Overall the *fcc* lattice is
24 essentially maintained on preparation. Upon thermal treatments structural and compositional
25 changes, in favour of alloyed structures with tetragonal packing, are seen.

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27 Keywords: X-ray Photoelectron Diffraction (XPD), Photoemission, Scanning Tunnelling
28 Microscopy (STM), Synchrotron, Surface Alloys

29 1. Introduction

30 Binary metal alloys have been the subject of much investigation,^{1, 2)} owing to their
31 widespread use in several fields spanning construction, transportation, electronics, medicine
32 and chemistry. Often the physical properties of bulk alloys are regarded as a linear
33 combination of the properties of the single components; in particular, the lattice parameter
34 is described by Vegard's law.³⁾ Copper-gold alloys find application in a variety of diverse
35 fields such as catalysis,⁴⁾ protection from corrosion⁵⁾ and jewellery.⁶⁾ The Cu/Au phase
36 diagram shows that copper and gold are completely miscible over the whole range of
37 concentrations and a large range of temperatures.^{2, 7, 8)} The alloy has a cubic unit cell, the
38 bulk lattice parameter of which positively deviates from Vegard's law.³⁾ Only a few ordered
39 phases are known:⁹⁾ for a one to one copper to gold stoichiometry, below *ca.* 700 K, two
40 ordered phases exist. The AuCu I phase, type *L1₀*, has a tetragonal superstructure and is
41 stable at low temperatures.¹⁰⁾ The second known stable phase is AuCu II, which differs from
42 AuCu I in that it has nearly periodic antiphase boundaries along the *y*-axis of the unit cell,
43 on average every five unit cells.¹⁰⁾ Other stable structures include Cu₃Au (type *L1₂*) and
44 CuAu₃ (type *L1₂*), in decreasing order of stability.⁷⁾ Recently, scanning tunnelling
45 microscopy (STM) investigations have focused on the surface layers of Cu/Au alloys,^{11, 12,}
46 ¹³⁾ prepared by annealing copper thin films deposited on Au(111) surfaces, highlighting the
47 growth mode and how ordered surface terminations can be prepared.¹³⁾ The nucleation
48 process of guest metals on Au(111) generally follows the place-exchange mechanism,¹⁴⁾
49 which was initially ruled out for the addition of copper; however, the same mechanism was
50 later confirmed also for copper.¹¹⁾

51 Although the nucleation step of guest metals on Au(111) seems now well understood,^{11,}
52 ^{12, 13, 14, 15, 16, 17, 18, 19, 20)} the morphology and composition of the top layer are still a matter of
53 debate. Specifically for the addition of copper, several studies report island growth without
54 copper intermixing with the underlying gold,^{21, 22, 23)} while others report intermixing for
55 every coverage even at room temperature.^{24, 25)} Complete encapsulation, whereby a single
56 copper layer is covered by a single gold layer, has also been proposed.^{12, 24, 25, 26)} Further, it
57 has been suggested that the structure and chemical composition of the layers below the first
58 play an important role in dictating the reactivity of such surfaces,^{5, 12, 13, 24, 26)} most likely due
59 to random intermixing between the two metals and chemically induced surface segregation
60 phenomena. In order to rationalize the properties of such structures, it is crucial to
61 characterize the layers just below the first, the top one having been successfully
62 characterized via STM.^{11, 12, 13)} This requires the use of techniques which can gather precise

63 information on a larger area than that evaluated via STM, typically a few square nanometres,
64 however, still maintaining sensitivity to the near surface layers.

65 In this contribution the top layers of structures prepared by depositing thin films of
66 copper on Au(111) are investigated using STM and synchrotron radiation based techniques:
67 X-ray photoelectron diffraction (XPD) allows one to access geometrical information on a
68 larger scale than STM, albeit averaged; and photoemission spectroscopy (PES) is used to
69 quantify elements and their chemical states, looking at both core level and valence band
70 regions. Both techniques are not only sensitive to the surface, but also to the layers
71 immediately below it.

72 Experimental results show structural and compositional changes upon annealing. The
73 determination of the structure of the uppermost layers is not trivial. However, through the
74 complementarity of the techniques employed, the overall geometry of the top layers is seen
75 to change from *fcc*, for multilayers of copper on Au(111), to tetragonal, when the near surface
76 alloys are formed.

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79 **2. Experimental methods**

80 The original STM investigation was carried out at the University of St. Andrews^{11, 13)}
81 and was used as a basis for the present work, which was undertaken at the PEARL
82 (Photoemission and Atomic Resolution Laboratory) beamline at the Swiss Light Source
83 (SLS).²⁷⁾ The beamline houses a variety of surface science techniques together with
84 synchrotron radiation methods allowing one to gather information on the surface and near-
85 surface layers. Base pressures in analysis, probe, and preparation chambers were all *ca.*
86 1×10^{-10} mbar.

87 The Au(111) crystal was cleaned by cycles of argon ion sputtering and annealing to *ca.*
88 870 K until wide terraces characteristic of the typical Au(111)- $(22 \times \sqrt{3})$ surface
89 reconstruction were observed via STM.^{28, 29)} Cleanliness after preparation of the Au(111)
90 crystal was verified by photoemission spectroscopy. The clean Au 4f core level spectra
91 showed the well-known low binding energy shifted surface component.³⁰⁾ Copper deposition
92 was achieved using a well outgassed FOCUS EFM evaporator and a 99.999% pure copper
93 rod. During copper deposition the pressure increased to *ca.* 8×10^{-10} mbar. The copper
94 deposition rate was *ca.* 0.02 ML min^{-1} . The coverage was calibrated collecting the Cu 2p_{3/2}
95 core level photoemission signal from a clean Cu(110) crystal as reference, at a photon energy
96 (PE) of 1045 eV, together with the evaluation of the attenuation of the Au 4f core level signal

97 recorded with a PE of 143 eV. Coverage evaluation is based on the assumption that the
 98 copper layer grows in a Frank-van der Merwe fashion³¹⁾ with no intermixing; a 1 ML
 99 coverage is defined as one copper atom per surface unit cell of Au(111). After preparation
 100 of the copper layer, the deposition was verified by photoemission spectroscopy. For surface
 101 sensitivity the PE for each of the core level transitions was chosen such as to obtain spectra
 102 at a kinetic energy of *ca.* 50 eV. The C 1s and O 1s spectra showed nothing above the noise,
 103 which, based on cross sections³²⁾ and escape depths³³⁾ at these kinetic energies, leads to upper
 104 limits of *ca.* 0.01 ML for both elements. Binding energy (BE) values are referred to the Fermi
 105 level collected at each PE. Various copper coverages were investigated from sub-monolayers
 106 to a few monolayers (*ca.* 2.7 ML). Analyser and beamline settings were chosen so that the
 107 detailed spectra were lifetime limited. The Cu 2p_{3/2} photoelectron diffraction data was
 108 collected over the complete 2π azimuth and a θ range up to 80° to the surface normal; details
 109 are reported in Ref.²⁷⁾ The photoelectron diffraction modulation^{34, 35, 36)} data were extracted
 110 using a 3×3 point averaging with a simple constant background subtraction and integration
 111 over the peak.

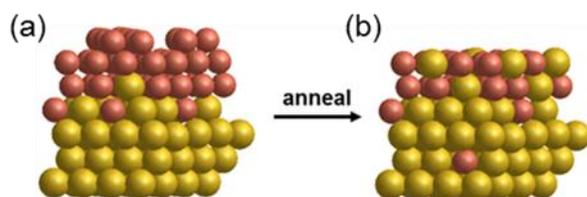
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114 3. Results and discussion

115 Following Ref.¹³⁾ the surface alloy was prepared by annealing a *ca.* 2.7 ML preparation
 116 to *ca.* 530 K. This process is represented schematically in figure 1. A copper layer *ca.* 2.7
 117 ML thick is deposited on Au(111), figure 1(a); some interdiffusion already occurs, as a result
 118 of the nucleation mechanism following the place-exchange mechanism.^{11, 13)} For coverages
 119 above 1 ML, the top layers show a moiré structure.^{12, 13, 26)} The system is then annealed to
 120 promote atomic mobility and diffusion, figure 1(b); as an effect of the annealing, copper
 121 sinks into the bulk gold. The extent of alloying is controlled by the annealing time and
 122 temperature.

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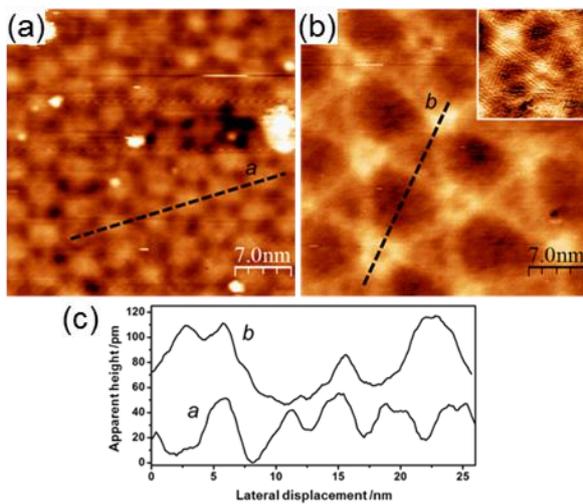
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125 **Fig. 1.** Schematics of the process for creating the surface alloy: (a) *ca.* 3 ML of copper
 126 (dark orange) are added to Au(111) (yellow); (b) the system is then annealed to promote
 127 atoms interdiffusion.

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In figure 2 STM images following preparation are reported. Figure 2(a) shows the moiré structure that forms for a copper coverage of 2 ML and above,^{12, 13, 26)} as a result of the interference between the added copper layers and Au(111) lattice. The appearance of a moiré pattern can be considered as an indication that only a small amount of intermixing between the two metals can occur. In principle the moiré pattern is generated by the lattice mismatch between the lattice unit cells of Cu(111), 0.256 nm, and Au(111), 0.288 nm. However, in the present case it is clear that considerable intermixing is occurring and that the moiré can also indicate the accommodation of the lattice space change from the surface layer (essentially copper) to the bulk (gold), going through a heavily alloyed interfacial region.



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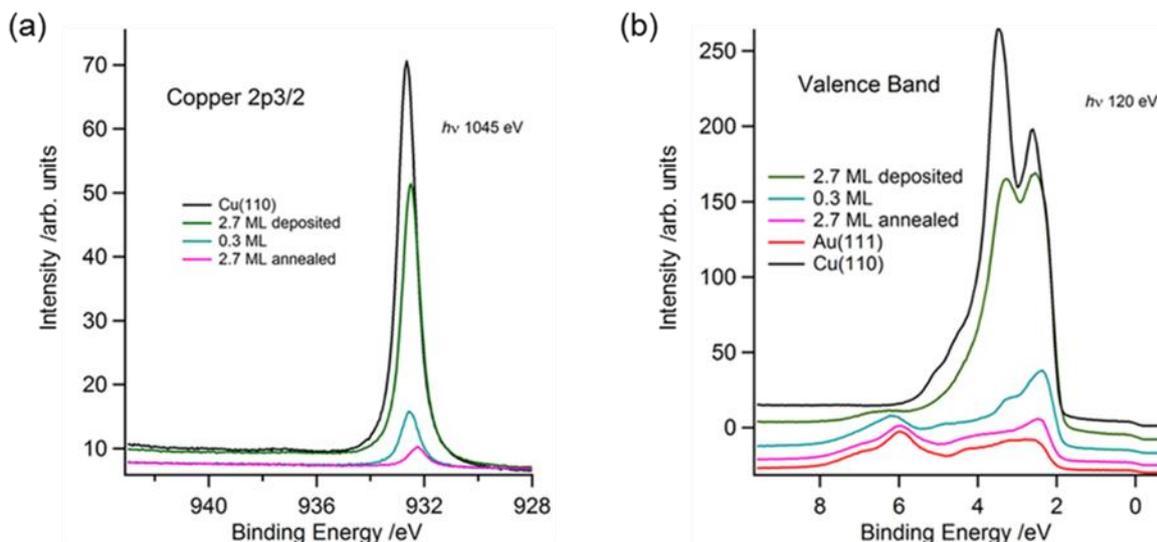
Fig. 2. STM images of the *ca.* 3 ML Cu/Au(111) before, (a), and after annealing, (b). (a) $35 \times 35 \text{ nm}^2$, 0.6 nA, -1.1 V; (b) $35 \times 35 \text{ nm}^2$, 0.7 nA, 1.1 V; inset $15 \times 15 \text{ nm}^2$, 0.3 nA, 0.4 V; (c) line profiles as in (a) and (b).

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For a coverage thicker than *ca.* 3 ML the system is reported to behave as a bulk copper {111} termination.^{12, 26)} The moiré has an average pseudo-periodicity of *ca.* 4.5 nm, as shown by the line profile *a* in figure 2(c). Upon annealing the surface alloy forms, figure 2(b); this is evidenced by surface features evolving into a different moiré structure having much larger, nearly double, average pseudo-periodicity, as shown by the line profile *b* in figure 2(c). The inset in figure 2(b) shows atomic resolution on the alloyed structure.

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Figure 3 shows Cu 2p_{3/2} core level and valence band photoemission spectra collected choosing the photon energies so that the kinetic energy was similar, so that a similar depth from the surface layer is probed, ensuring the same surface sensitivity.



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156 **Fig. 3.** Cu 2p3/2 photoemission (a) and valence band (b) spectra collected at copper
 157 increasing coverage (before annealing) and after annealing. Valence band spectra have been
 158 offset for clarity.

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160 Figure 3(a) shows the Cu 2p3/2 core level spectra for three preparations together with
 161 that of a Cu(110) crystal. The quoted coverages are based on a simple quantification based
 162 on a layer-by-layer growth mode which, considering the STM observations^{12, 13, 26)} is clearly
 163 inappropriate to describe the growth process in full. Nevertheless, given the complexity of
 164 both nucleation and growth steps, this is a simple and prudent choice. On closer observation,
 165 there is a shift in the binding energy of the Cu 2p3/2 spectra from 932.1 eV at the lowest
 166 submonolayer coverage, to that of bulk copper, as measured on a clean Cu(110) single crystal,
 167 at 932.5 eV. At *ca.* 3 ML the BE value (932.4 eV) is very close to that of bulk copper. The
 168 overall peak shift across the range of coverages investigated is 0.4 eV, a much smaller value
 169 than its FWHM, *ca.* 0.85 eV, which remains constant. STM investigations^{11, 12, 13, 26)} show
 170 that the morphology of the top layers undergoes considerable changes and cannot be
 171 associated with the presence of a single metal. In particular for the low coverage regime
 172 (below *ca.* 1 ML), this indicates that there is a subsurface region which is rather rich in
 173 copper. The presence of a subsurface copper-rich layer for the Cu/Au system is in agreement
 174 with several computational studies.^{37, 38, 39, 40)} The layers above the first exhibit a moiré
 175 structure instead; also in these layers some intermixing is seen,^{12, 26)} however to a lesser
 176 extent than for the interfacial layer. At above *ca.* 3-4 ML, copper thin films exhibit electronic
 177 characteristics similar to those of bulk copper.^{12, 26)} This can be explained on the basis of the
 178 relative mobility of the atoms in the uppermost surface layers, when compared to that of the
 179 bulk-like layers underneath. As a consequence of such a variable degree of intermixing,
 180 different growth modes are experienced. In this case one might expect various copper species

181 to be identified by distinct binding energy shifts; however this is not observed. Such a shift
182 was not observed by Zhao and co-workers,²⁴⁾ likely because only the dissolution of 1 ML
183 thick copper film into bulk gold was investigated. In such a system, the copper layer is sub-
184 surface already after preparation and the different chemical environments that copper atoms
185 experience when dissolving further into the bulk may be difficult to observe because of
186 differences in escape depths and also the fact that a conventional non-monochromatic X-ray
187 source was used compared to a monochromatized synchrotron source as at PEARL.

188 On the basis of previous experimental work, ion scattering and X-Ray Photoelectron
189 Spectroscopy (XPS)²⁴⁾ and STM,^{11, 12, 13, 26, 41)} it is clear that the copper is easily incorporated
190 below the gold surface at room temperature even at sub-monolayer coverage. At higher
191 copper coverages the two metals interdiffuse forming a copper rich subsurface region. The
192 gold concentration in the copper layers is seen to decrease with increasing copper coverage
193 (the thicker the copper layer, the lower the gold content). The diffusion of copper into the
194 bulk gold is facilitated with increasing temperature. This is seen for the 2.7 ML coverage
195 which, on annealing, decreases to a value which nominally corresponds to a *ca.* 0.1 ML,
196 clearly demonstrating the dissolution of the copper into the gold bulk.

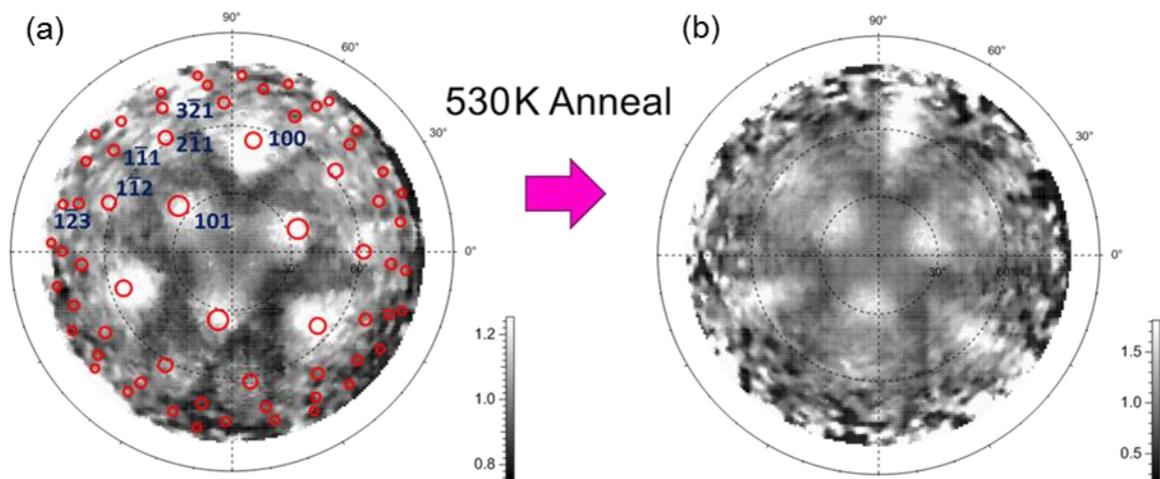
197 Figure 3(b) shows valence band measurements recorded for the same copper coverages
198 as in figure 3(a) and that of a clean Au(111) sample. Here of interest is the relative signal
199 intensity, when compared to that of the Cu 2p_{3/2} core level, and also the intensity of the
200 copper and gold contributions to the valence band. The latter can be understood in terms of
201 the relative size of the cross sections of the Cu 3d and Au 5d levels, 5.139 Mb and 0.4332
202 Mb at a PE of 132.3 eV, respectively.³²⁾ The very intense valence band signal of the 2.7 ML
203 copper preparation decreases by almost a factor of 50 on annealing; by comparison, the
204 decrease in intensity of the Cu 2p_{3/2} is about a factor of 15. This is difficult to explain as the
205 kinetic energies of the electrons are very similar (Cu 2p_{3/2} 113 eV, valence band 117 eV) so
206 the depth probed by the photoelectrons is similar. The reason for the different decrease in
207 signal is probably due to alloying. The electronic structure of the bonding orbitals will
208 change with alloy composition and consequently the valence band photoemission will
209 change. Indeed, in all coverage regimes up to and including the 2.7 ML simple weighted
210 sums are unable to explain the changes in the valence band, indicating complex growth is
211 present over the whole range studied. This is also supported by the observed shifts in the Cu
212 2p_{3/2} core levels.

213 As discussed above, the growth and dissolution of copper in the Au(111) surface are
214 complex processes and, as such, may occur following different mechanisms. However, with

215 a controlled annealing process, near surface alloy with an ordered structure can be formed.
 216 This is shown in figure 2(b) which shows the moiré that is formed on annealing.¹³⁾ A short
 217 range ordered structure is seen within the moiré, as highlighted in the inset in figure 2(b).
 218 The question now arises as to whether this is purely a local ordering on the uppermost surface
 219 layers or is connected to the bulk; in the latter case, what is the relation between the surface
 220 layer and the bulk. This is a difficult and complicated question and to gain an insight into
 221 this the synchrotron related technique of angle resolved XPD was chosen as both surface
 222 and signals from buried layers can be readily distinguished. Photoelectron diffraction is a
 223 well-established laboratory and synchrotron technique for structural analysis.^{34, 35, 36, 42)} It is
 224 also chemical element and state sensitive from the choice of the photoemission peak. Any
 225 disorder will be manifested by the absence of structure in the diffraction pattern.

226 Cu 2p_{3/2} XPD patterns for a PE of 1045 eV from the 2.7 ML preparation before and after
 227 annealing are shown in Figure 4. In both diffraction patterns, structure is seen. An increase
 228 in intensity modulation is seen from before, figure 4(a) to after annealing, figure 4(b).

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231 **Fig. 4.** XPD patterns of Cu 2p_{3/2} (PE 1045 eV) for the 2.7 ML preparation before (a) and
 232 after annealing (b). Crystallographic directions of an *fcc* lattice have been superposed to the
 233 2.7 ML pattern (a); see text for further details. An increased ordering as a result of annealing
 234 can be seen in the range bar of pattern (b).

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238 The patterns are very different from each other, thus indicating clear changes in the
 239 geometrical structure. Superposed on figure 4(a) is a pattern obtained from the stereographic
 240 projection of the *fcc* lattice, where the forward scattering peaks of the different scattering
 241 directions are depicted by circles and their intensities by the circle size. Such an approach
 242 can be regarded as a first approximation to a single-scattering calculation.⁴³⁾ The model

243 reproduces the experimental pattern fairly well, with scattering in the $\langle 110 \rangle$, $\langle 100 \rangle$ and
244 $\langle \bar{1}12 \rangle$ directions; higher orders spots are considerably weaker or absent in the data, when
245 compared to the model. Based on other literature studies^{12, 13, 24)} this suggests that the copper
246 grows already as an alloy, or alloys, in an *fcc* structure with some disorder extending into the
247 bulk. After annealing the pattern changes considerably, figure 4(b), however retains the
248 principal $\{111\}$ orientation of the underlying gold surface, and the main peaks grow in
249 intensity. There are some small shifts in the $\langle 110 \rangle$ and $\langle 100 \rangle$ directions, the
250 $\langle \bar{1}12 \rangle$ essentially disappear, and a relatively intense spot appears in the $\langle 22\bar{1} \rangle$ directions.
251 The disappearance of the $\langle \bar{1}12 \rangle$ spots can be due to a lowering of translational symmetry
252 which is consistent with the geometry of the surface layers changing from an *fcc* lattice (Cu)
253 to tetragonal (AuCu). The changes seen in the XPD pattern after annealing, figure 4(b),
254 require further investigation, and modelling. The data here presented is dominated by
255 forward scattering with little visible fine, secondary, structure. Perhaps this is not surprising,
256 given the formation of a reconstruction with an averaged pseudo-periodicity larger than that
257 of the copper film initially deposited, and some variation in the local order, as shown in
258 figure 2(b). Consequently, the primary information obtained is an averaged lattice symmetry,
259 whilst limited geometrical length information can be extracted from triangulation.
260 Nevertheless, it is clear that a near surface alloy forms, the symmetry of which is no longer
261 that of an ideal *fcc* lattice, but consistent with a distortion of the *fcc* lattice to a lower
262 symmetry, such as a tetragonal $L1_0$. More detailed modelling of the tetragonal AuCu alloys,
263 the related AuCu I and AuCu II phases, and other structures is being investigated.

264 Both observations, the growth in a preferred orientation and essentially its retention, and
265 the structural change after annealing, given the large area of the surface covered by the moiré
266 and other structures seen in the STM experiments, indicate that the underlying Au(111)
267 surface is still able to influence the growth of the alloys and the interfaces.

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269 **4. Conclusions**

270 A combined synchrotron and STM study on the growth of copper layers on the Au(111)
271 surface and alloying on annealing has been undertaken. Preliminary data analysis show
272 structural changes associated with the alloying process. Upon preparation, copper
273 multilayers (*ca.* 3 ML) exhibit an overall *fcc* geometry, as determined by XPD. However,
274 STM shows the appearance of a moiré pattern, which is ascribed to the interference of the
275 copper and Au(111) lattices. On annealing the systems evolves into a near-surface alloy;

276 XPD reveals a structural change, in favour of a tetragonal geometry, and STM shows very
277 different topographical features.

278 Geometrical optimization of model structures based on the experimental data and the
279 resulting electronic structure are being pursued in order to obtain a much improved
280 description, in the attempt to rationalize the interdiffusion of the atoms and predict the
281 reactivity of such a system.

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291

292 **Appendix**

293 Additional experimental details: Survey scans using a PE of 1250 eV (for comparison with
294 laboratory Mg K α work) were taken along with detailed scans of the various core levels of
295 interest and valence region. For a comparison with bulk like spectra Cu 2p, Au 4f and valence
296 band spectra were taken with a PE of 1400 eV.

297 STM experimental: In-house STM measurements were performed on an Omicron variable
298 temperature STM, operated at room temperature, in an UHV environment with a base
299 pressure of about 1×10^{-10} mbar or lower. Topographs were acquired at room temperature
300 in constant current mode using electrochemically etched tungsten tips. The tunneling voltage
301 bias was applied to the tip, whereas the sample was grounded. Copper was deposited on the
302 Au(111) crystal kept at room temperature by electrically heating a high purity copper wire
303 (99.999%, 0.1 mm diameter), wrapped around a tantalum wire (99.999% purity, 0.25 mm
304 diameter), to yield a deposition rate of *ca.* 0.07 ML min⁻¹. The nominal coverage in
305 monolayer (ML) was determined by evaluating the fraction of each image covered by the
306 features related to copper and their appearance, in combination with the calculated exposure.

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Figure Captions

Fig. 1. Schematics of the process for creating the surface alloy: (a) *ca.* 3 ML of copper (dark orange) are added to Au(111) (yellow); (b) the system is then annealed to promote atoms interdiffusion.

Fig. 2. STM images of the *ca.* 3 ML Cu/Au(111) before, (a), and after annealing, (b). (a) $35 \times 35 \text{ nm}^2$, 0.6 nA, -1.1 V; (b) $35 \times 35 \text{ nm}^2$, 0.7 nA, 1.1 V; inset $15 \times 15 \text{ nm}^2$, 0.3 nA, 0.4 V; (c) line profiles as in (a) and (b).

Fig. 3. Cu 2p_{3/2} photoemission (a) and valence band (b) spectra collected at copper increasing coverage (before annealing) and after annealing. Valence band spectra have been offset for clarity.

Fig. 4. XPD patterns of Cu 2p_{3/2} (PE 1045 eV) for the 2.7 ML preparation before (a) and after annealing (b). Crystallographic directions of an *fcc* lattice have been superposed to the 2.7 ML pattern (a); see text for further details. An increased ordering as a result of annealing can be seen in the range bar of pattern (b).

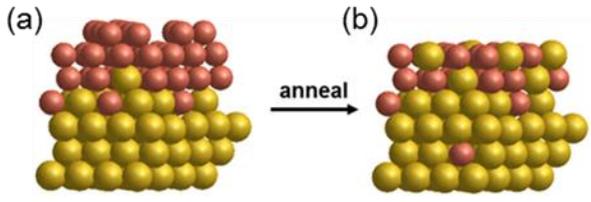


Fig.1.

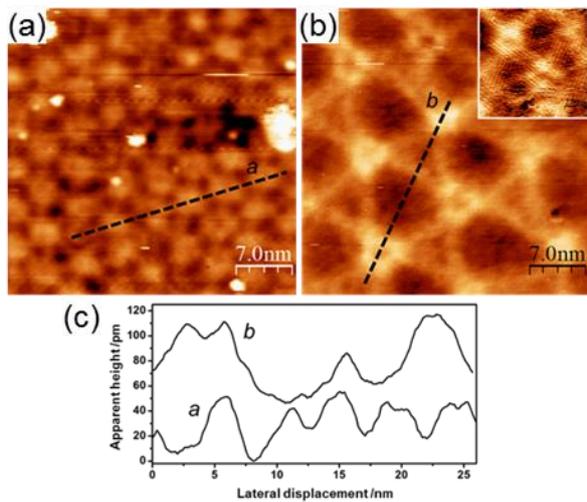


Fig.2.

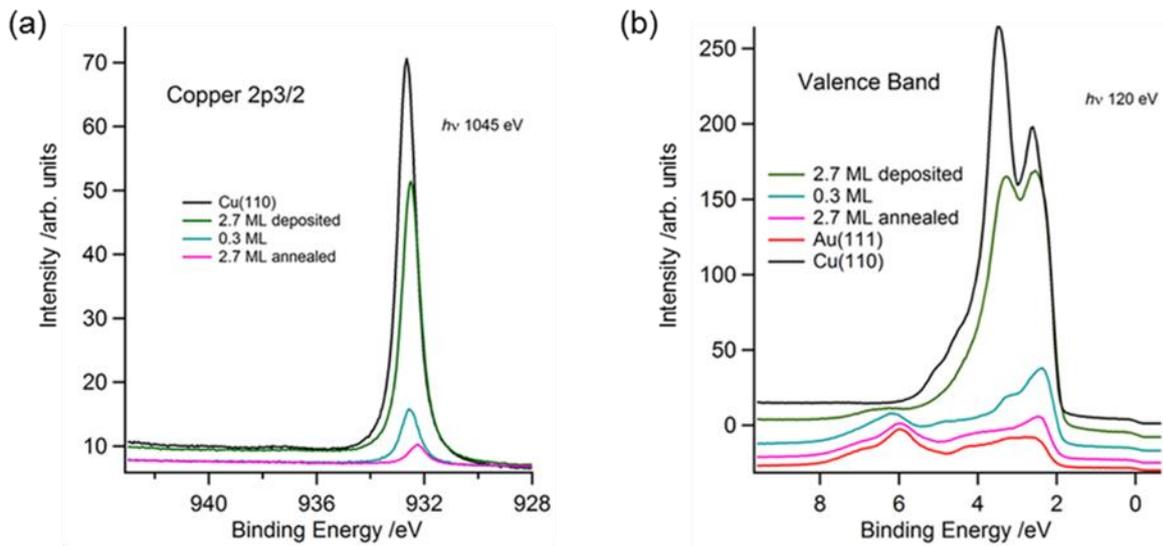


Fig.3.

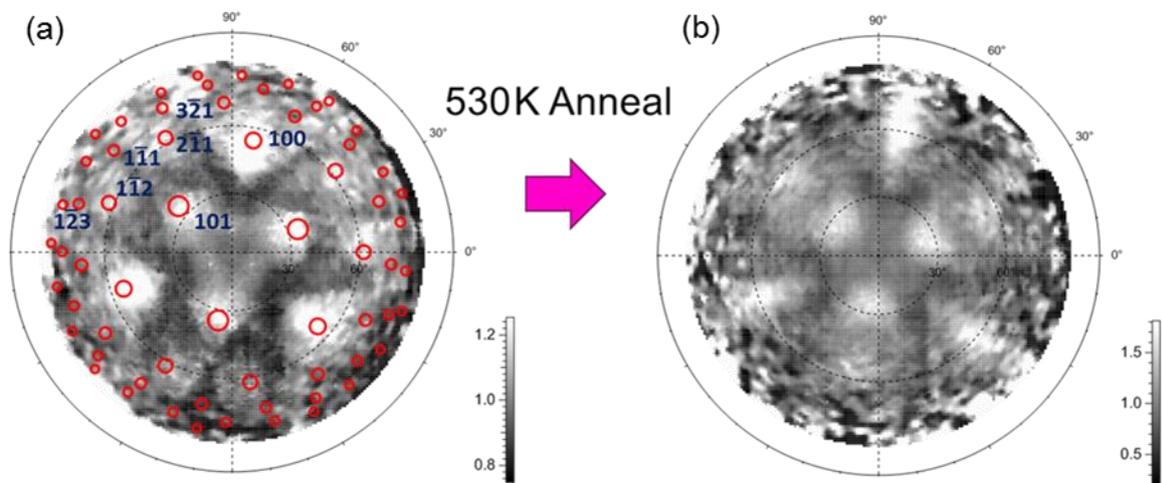


Fig.4.