

# Planets around evolved intermediate-mass stars in open clusters<sup>★</sup>

## II. Are there really planets around IC4651No9122, NGC2423No3 and NGC4349No127?

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### Abstract

**Aims.** The aim of this work is to search for planets around intermediate-mass stars in open clusters by using data from an extensive survey with more than 15 years of observations.

**Methods.** We obtain high-precision radial velocities (*RV*) with the HARPS spectrograph for a sample of 142 giant stars in 17 open clusters. We fit Keplerian orbits when a significant periodic signal is detected. We also study the variation of stellar activity indicators and line-profile variations to discard stellar-induced signals.

**Results.** We present the discovery of a periodic *RV* signal compatible with the presence of a planet candidate in the 1.15 Gyr open cluster IC4651 orbiting the 2.06 M<sub>⊙</sub> star No. 9122. If confirmed, the planet candidate would have a minimum mass of 7.2 M<sub>J</sub> and a period of 747 days. However, we also find that the FWHM of the CCF varies with a period close to the *RV*, casting doubts on the planetary nature of the signal. We also provide refined parameters for the previously discovered planet around NGC2423 No. 3 but show evidence that the BIS of the CCF is correlated with the *RV* during some of the observing periods. This fact advises us that this might not be a real planet and that the *RV* variations could be caused by stellar activity and/or pulsations. Finally, we show that the previously reported signal by a brown dwarf around NGC4349 No. 127 is presumably produced by stellar activity modulation.

**Conclusions.** The long-term monitoring of several red giants in open clusters has allowed us to find periodic *RV* variations in several stars. However, we also show that the follow-up of this kind of stars should last more than one orbital period to detect long-term signals of stellar origin. This work warns that although it is possible to detect planets around red giants, large-amplitude, long-period *RV* modulations do exist in such stars that can mimic the presence of an orbiting planetary body. Therefore, we need to better understand how such *RV* modulations behave as stars evolve along the Red Giant Branch and perform a detailed study of all the possible stellar-induced signals (e.g. spots, pulsations, granulation) to comprehend the origin of *RV* variations.

**Key words.** stars: individual: IC4651 No. 9122, NGC2423 No. 3, NGC4349 No. 127 – stars: planetary systems – stars: evolution – planets and satellites: detection – Galaxy: open clusters and associations

### 1. Introduction

In the last 20 years more than 3500 planets have been discovered, mainly around Main Sequence (MS) solar type stars (exoplanet.eu, Schneider et al. 2011). One of the most successful methods to detect planets, the radial velocity (*RV*) technique, is more difficult to be used in stars hotter than ~6500 K due to the increase in rotational velocities of those stars and the lack of a sufficient number of spectral lines to determine the velocity shifts in their spectra. Nevertheless, several brown dwarfs and planets have been reported by adapting the *RV* technique to A-

F stars (Galland et al. 2006; Desort et al. 2008; Lagrange et al. 2009; Borgniet et al. 2014, 2016). However, in order to understand the planetary formation mechanisms around early F or A stars we need larger samples, together with the determination of planetary masses. Thus, the preferred option to solve this issue has been to apply the *RV* method on K giants, the evolved counterparts of those massive stars, with generally low rotation rates and a larger number of spectral lines in their much cooler spectra when compared with their unevolved MS counterparts (e.g. Frink et al. 2002; Sato et al. 2003; Lovis & Mayor 2007; Niedzielski et al. 2015). This allows to probe a different stellar mass range.

A major issue when interpreting *RV* variations in red giants is the presence of intrinsic stellar jitter which shows a typical

<sup>★</sup> Based on observations collected at the La Silla Observatory, ESO (Chile), with HARPS/3.6m (runs ID 075.C-0140, 076.C-0429, 078.C-0133, 079.C-0329, 080.C-0071, 081.C-0119, 082.C-0333, 083.C-0413, 091.C-0438, 092.C-0282, 099.C-0304 and 0100.C-0888) and with UVES/VLT at the Cerro Paranal Observatory (run 079.C-0131)

**Table 1.** Stellar characteristics of the analyzed planet host candidates. Distances and magnitudes are extracted from WEBDA and SIMBAD, respectively.

		IC4651 No. 9122	NGC2423 No. 3	NGC4349 No. 127
$T_{\text{eff}}$	K	$4720 \pm 71$	$4592 \pm 72$	$4503 \pm 70$
$\log g$	( $\text{cm s}^{-2}$ )	$2.72 \pm 0.15$	$2.33 \pm 0.20$	$1.99 \pm 0.19$
$[Fe/H]$		$0.08 \pm 0.04$	$-0.03 \pm 0.05$	$-0.13 \pm 0.04$
$M$	$M_{\odot}$	$2.06 \pm 0.09$	$2.26 \pm 0.07$	$3.81 \pm 0.23$
$R$	$R_{\odot}$	$8.90 \pm 0.68$	$13.01 \pm 1.11$	$36.98 \pm 4.89$
$L$	$L_{\odot}$	35.29	67.56	504.36
Age	Gyr	$1.223 \pm 0.150$	$0.850 \pm 0.068$	$0.203 \pm 0.031$
Distance	pc	888	766	2176
$V$	mag	10.91	$10.04 \pm 0.04$	$10.82 \pm 0.08$
$\alpha$		17:24:50.1	07:37:09.2	12:24:35.5
$\delta$		-49:56:56.1	-13:54:24.0	-61:49:11.7

level of  $10\text{-}15 \text{ m s}^{-1}$  (e.g. Setiawan et al. 2004; Hatzes et al. 2005) and increases towards more evolved stages (Hekker et al. 2008) and for redder stars (Frink et al. 2001). These short-term (few hours to few days) low-amplitude  $RV$  variations are mainly caused by radial oscillations, i.e. p-modes (e.g. Hatzes et al. 2007). On the other hand, long-term non-radial pulsations can produce larger  $RV$  amplitudes of hundreds of  $\text{m s}^{-1}$  (e.g. Hatzes & Cochran 1999) which eventually can mimic the presence of a planet. Moreover, the combination of several modes of radial oscillations can produce night-to-night variations of  $\sim 100 \text{ m s}^{-1}$  (e.g. see the case of  $\alpha$  Boo and  $\alpha$  Tau, Hatzes & Cochran 1993). The modulation of active regions in red giants can produce large amplitude  $RV$  variations and on longer timescales as well. Therefore, it is important to carry out long-term observations covering more than one period of the planet candidates (and for a time span larger than the stellar rotational period), to evaluate the stability of the hypothetical planetary signal and its possible relation with the rotational period of the star.

Some interesting correlations between metallicity, stellar mass and the presence of planets have been proposed as more planets have been discovered during the last years. Johnson et al. (2010) found that the frequency of massive planets is higher around more massive stars but later, Reffert et al. (2015) reported that the giant planet occurrence rate peaks at  $\sim 1.9 M_{\odot}$  and then rapidly drops for masses larger than  $\sim 2.5\text{-}3 M_{\odot}$ . They also found a clear planet-metallicity correlation for their confirmed planet-host giant star sample. However, the works by Maldonado et al. (2013) and Mortier et al. (2013) only found such correlation for giant stars with masses larger than  $\sim 1.5 M_{\odot}$ .

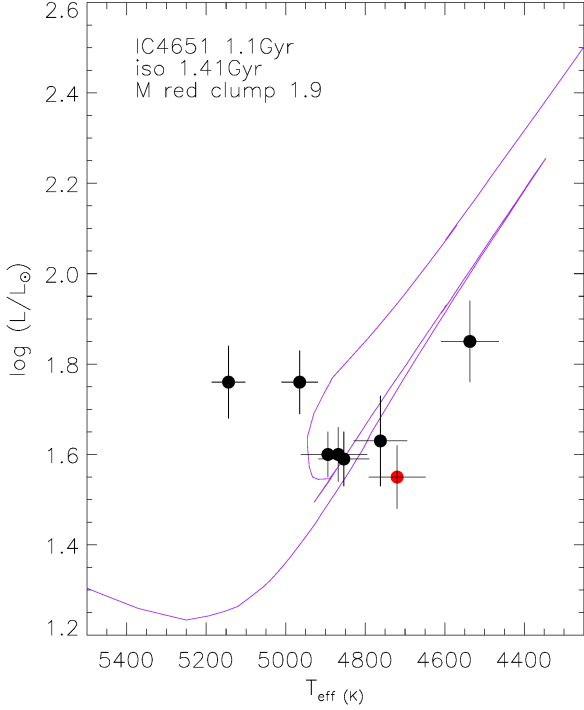
Therefore, it is clear that in order to understand planet formation mechanisms we need to get accurate masses for the host stars. One of the main problems in characterizing these *evolved* planetary systems is that the determination of masses for red giants is complicated by the crowding of evolutionary tracks with similar stellar parameters in that region of the Hertzsprung-Russell (HR) diagram and thus the determined planetary masses with the  $RV$  technique may present very large uncertainties (Lloyd 2011; Sousa et al. 2015) (but see Ghezzi & Johnson 2015; North et al. 2017). As a way to address this issue several programs to discover planets in open clusters have been started. The advantage of working with clusters is that ages and masses of their stars can be better constrained, thus the planetary characterization will be more reliable. Despite the efforts to detect those planets, only few discoveries have been reported. The first planet ever discovered in an open cluster was announced by Sato et al. (2007). This was a long period planet (595 days)

with a minimum mass of  $7.6 M_J$  orbiting a red giant of  $2.7 M_{\odot}$  in the Hyades ( $\sim 600$  Myr). Soon after, a similar planet was announced around a  $2.4 M_{\odot}$  giant in NGC2423 ( $m_2 \sin i = 10.6 M_J$ ,  $P = 750$  days) by the  $RV$  survey around intermediate mass stars of Lovis & Mayor (2007, hereafter Paper I). This work also reported the discovery of a brown dwarf orbiting a  $3.9 M_{\odot}$  red giant in NGC4349. Finally, Brucalassi et al. (2014) discovered a Jupiter-like planet in a 120 days orbit around a K giant in M67. We note that the search of planets in open clusters has also targeted MS stars with some detections of Neptune-size planets (Meibom et al. 2013) and hot Jupiters in different clusters (Quinn et al. 2012, 2014; Malavolta et al. 2016; Brucalassi et al. 2014, 2016).

The aim of this work is to present the new results of our  $RV$  survey started in Paper I which includes the discovery of a long-period planet candidate in IC4651. The outline of the paper is as follows: in Sect. 2 we present the data and derivation of stellar parameters. The planet candidate found in IC4651 is debated in Sect. 3 and the discussion of the signals previously attributed to the presence of a planet in NGC2423 is presented in Sect. 4. In Sect. 5 we provide evidence that the previously discovered signal from a brown dwarf in NGC4349 is most likely caused by modulation of stellar magnetic activity. Finally, in Sects. 6 and 7 we present a general discussion and conclusions of the results.

## 2. Observations and sample

The  $RV$  sample used in this work is fully described in Paper I. The objects analyzed here have been followed for nearly 5 years (from March 2005 to October 2009) with HARPS (ESO-3.6m, La Silla) and some of them were also observed with CORALIE (1.2 m-Swiss Telescope, La Silla) in previous years. In summary, the survey has focused in open clusters observable from La Silla, with bright giants (maximum  $V = 10$  or  $13$ , for each instrument respectively) having masses between  $1.5$  and  $4 M_{\odot}$ . Moreover, only clusters with at least 3 giants known to be non-binary cluster members were chosen. Additional observations have been taken during 2017 and 2018 for the targets showing large  $RV$  variations. The observing programs are detailed in a footnote in page 1. The observations were done using classical fibre spectroscopy mode (no simultaneous calibration) and the exposures times were estimated in order to have individual spectra with a signal-to-noise ratio (S/N)  $\sim 30$ . This gives a typical  $RV$  photon-noise of  $\sim 3.5 \text{ m s}^{-1}$  which is enough to detect massive planets around the surveyed stars.



**Figure 1.** HR diagram for the cluster IC4651. The purple line is the isochrone of 1.25 Gyr and  $Z = 0.018$  ( $[Fe/H] = 0.09$  dex). The red dot is the planet host IC4651 No. 9122.

In total, 142 stars were monitored within 17 open clusters using the HARPS spectrograph at the ESO 3.6m telescope ( $R \sim 115000$ ). The stellar parameters, namely the effective temperature ( $T_{\text{eff}}$ ), surface gravity ( $\log g$ ), metallicity ( $[Fe/H]$ ), and microturbulence ( $\xi_r$ ) for most of the clusters were presented in Santos et al. (2009, 2012) and improved upon in Delgado Mena et al. (2016) (by using higher S/N data and a Fe linelist optimized for cool stars) together with the derivation of stellar ages, masses and radii. In this paper we will focus on the results for three stars in the clusters IC4651, NGC2423 and NGC4349. The parameters for those targets are summarized in Table 1.

### 3. A planet candidate around IC4651 No. 9122

#### 3.1. Parent star characteristics

Our sample contains 8 giant stars in the open cluster IC4651 (distance  $d = 888$  pc) with an average metallicity  $[Fe/H] = 0.06 \pm 0.06$  dex (see Table B.1 of Delgado Mena et al. 2016). In Fig. 1 we show the HR diagram for this cluster. We use PARSEC isochrones (Bressan et al. 2012) to derive masses and ages (see Table B.2 of Delgado Mena et al. 2016). The isochrone which better matches our stars has an age of 1.41 Gyr, slightly higher than the 1.14 Gyr provided by the WEBDA database<sup>1</sup>. Our planet candidate host, with  $T_{\text{eff}} = 4720 \pm 71$  K,  $\log g = 2.72 \pm 0.15$  dex,  $M = 2.06 \pm 0.09 M_{\odot}$  and  $R = 8.90 \pm 0.68 R_{\odot}$  (see Table 1) seems to be on the first ascent of the Red Giant Branch (RGB) while most of the other stars seem to be located in the red clump. The mean radial velocity of the giants in this cluster is  $-30.28 \pm 1.33$  km s<sup>-1</sup> while the mean RV of IC4651 No. 9122 is  $-30.26$  km s<sup>-1</sup>, thus this star is likely a cluster member. We estimated a  $v \sin i$  of  $3.89$  km s<sup>-1</sup> in

Delgado Mena et al. (2016) for a fixed macroturbulence velocity of  $2.36$  km s<sup>-1</sup> (with the empirical formula by Valenti & Fischer 2005). This leads to a maximum rotational period of  $\sim 115$  days. If we consider the empirical relations for macroturbulence given by Gray (2005) and Hekker & Meléndez (2007) for class III stars we would obtain a macroturbulence velocity of  $5.25$  km s<sup>-1</sup> which would then cause  $v \sin i$  to decrease to  $0.1$  km s<sup>-1</sup> (the lowest limit accepted by the procedure to derive  $v \sin i$ ). This is a probably unrealistic value which would lead to a rotational period of 4503 days.

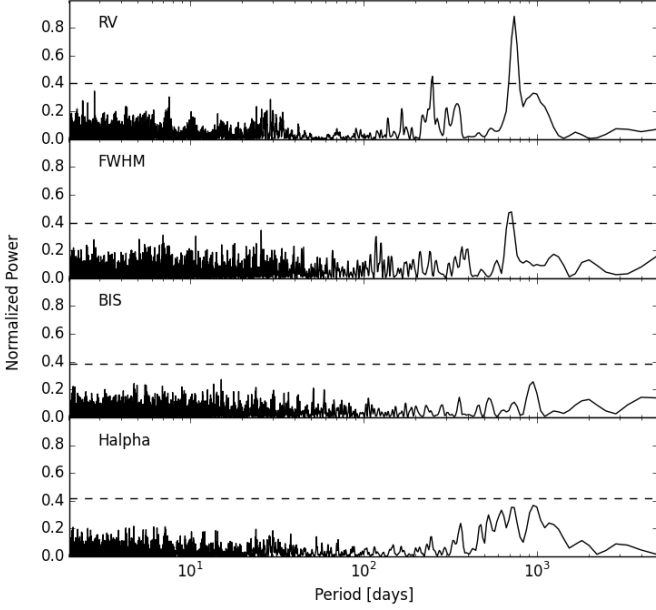
In order to illustrate the stellar jitter of the observations we plot the RV variations of the stars in the cluster in Fig. A.1. Most of the stars exhibit a scatter around  $13$  m s<sup>-1</sup> while our planet candidate host shows a dispersion of  $60$  m s<sup>-1</sup>, clearly above the average stellar jitter. Thus, this may hint to the presence of a planet. The star with the RV dispersion of  $26$  m s<sup>-1</sup> is IC4651 No. 9791, which show RV peak to peak variations of  $80$  m s<sup>-1</sup>. The photometric survey by Sahay et al. (2014) reported that this star is a long-period variable but could not establish its period due to the limited duration of the observations (130 days). Therefore, it is probable that the variability observed in brightness is related with the scatter in RV and of stellar origin, although it is not possible either to determine any period with our RV data. Interestingly, other two stars in this cluster, No. 8540 and No. 9025 also show brightness variability as reported by Sahay et al. (2014), which in the case of the latter reaches 0.4 magnitudes with a 73 days cadence. However, the RV scatter for these two stars is lower,  $10.8$  m s<sup>-1</sup> and  $12.1$  m s<sup>-1</sup>, respectively, of the level of the expected noise in these stars. To diagnose these effects we will analyse different activity proxies in Section 3.4.

#### 3.2. The radial velocity dataset

In the upper panel of Fig. A.2, the RV data spanning 12.5 years of observations is shown. We collected a total of 47 points between ESO periods 75 and 83. These data already allow to infer the possible presence of a planet but we found hints that the RV variability might be of stellar origin. Therefore, we collected 4 RV additional points<sup>2</sup> during periods 99 and 100 (between April and October 2017). Moreover we found in the ESO archive 6 extra points taken during 2013 and 2014 (programs 091.C-0438 and 092.C-0282 by Sanzia Alves et al.) as listed in Table B. The typical photon noise dominated errors in RV of the data is  $\sim 3$  m s<sup>-1</sup>. However, the RV variations are usually dominated by the stellar jitter which is much higher in this kind of stars as discussed in previous subsection. For example, solar-like p-mode oscillations in this star are expected to produce an amplitude of  $\sim 4$  m s<sup>-1</sup> and a period of 0.13 days estimated with the scaling relations of Kjeldsen & Bedding (1995). Since the average scatter in the RV data (for the 'stable' stars) is close to  $15$  m s<sup>-1</sup> we decided to add quadratically this noise to the photon noise before fitting the data to determine the planet orbit. We used the *yorbit* algorithm Ségransan et al. (2011) to fit the whole dataset with a model composed of a Keplerian function. *Yorbit* uses a hybrid method based on a fast linear algorithm (Levenberg-Marquardt) and genetic operators (breeding, mutations, crossover), and has been optimised to explore the parameter space for Keplerian fit-

<sup>2</sup> We note that we have applied a small negative offset to the RV points taken after 2015, when an upgrade of the HARPS fibers took place (Lo Curto et al. 2015). The shift was calculated by extrapolating the measurements of Table 3 in Lo Curto et al. (2015) to the FWHM of the spectra for this star.

<sup>1</sup> <http://www.univie.ac.at/webda/>



**Figure 2.** Generalized Lomb-Scargle periodograms of RV, FWHM, BIS and H $\alpha$  index for IC4651 No. 9122. The dashed line indicates the FAP at 1% level.

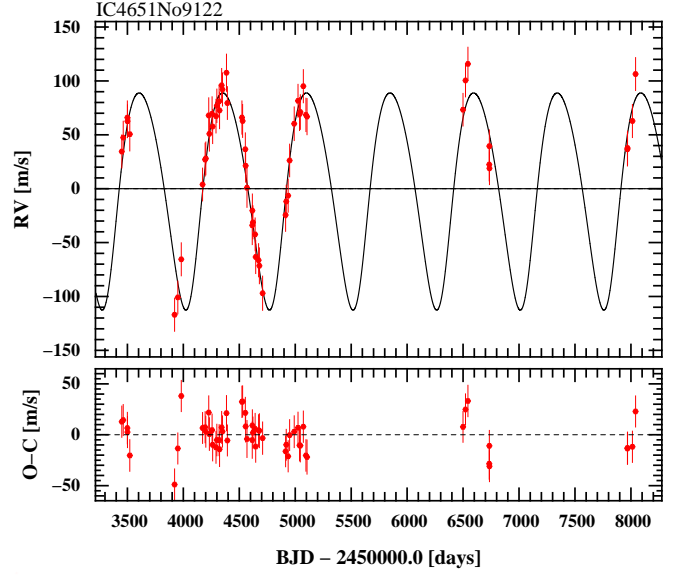
ting of radial velocity datasets. In our case, the global search for orbital parameters was made with a genetic algorithm.

An analysis using Generalised Lomb-Scargle (GLS) periodograms (e.g. Zechmeister & Kürster 2009) was performed for RV (see upper panel of Fig. 2). The false-alarm-probability (FAP) was computed by bootstrapping the data and statistically significant peaks were considered for values above FAP = 1%. A significant signal with a period of  $\sim 741$  days can be clearly observed in the periodogram. This periodic variation is well fitted by a Keplerian function with  $P = 747$  days,  $K = 100.8 \text{ m s}^{-1}$ , and  $e = 0.15$  (see Table 2). These values correspond to the expected signal induced by a planet with  $7.2 M_J$  in a 2.05 AU semi-major orbit. The phase curve of the best fitted solution can be seen in Fig. 3. The dispersion of the residuals is  $17 \text{ m s}^{-1}$  and the reduced  $\chi^2$  is 1.34. As a further test we run a Markov chain Monte Carlo (MCMC) to evaluate the white noise in our data. We get orbital parameters similar to those found by *yorbit* and we determine a jitter of  $13.5 \text{ m s}^{-1}$  which is very close to the average RV in the cluster. Therefore, we can conclude that the assumption of adding a noise of  $15 \text{ m s}^{-1}$  does not alter the results and it is an acceptable solution to treat the stellar-induced noise that cannot be properly modelled with our current data.

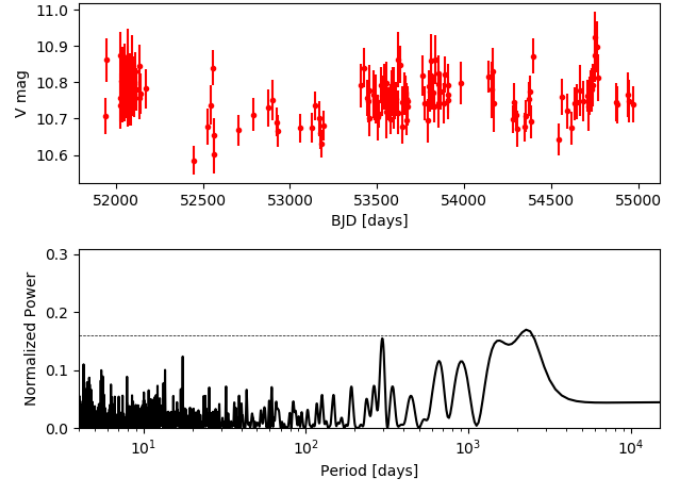
### 3.3. Photometry

This star has been observed within the All Sky Automated Survey (ASAS) from the Las Campanas Observatory (Chile) with observations available between 2001 and 2009 (Pojmanski 2002), hence contemporaneous to our first set of RV measurements. We found 151 measurements classified as good quality (they are given the grade A or B, with average errors of 0.05 mag) which show a  $\sim 0.3$  peak-to-peak variability in V magnitude. We downloaded the light curves in V magnitude from the The ASAS-3 Photometric V-band Catalogue<sup>3</sup> and performed a GLS to detect any possible variability with the same period as

<sup>3</sup> <http://www.astrouw.edu.pl/asas/?page=aasc&catsrc=asas3>



**Figure 3.** Radial velocity curve as a function of time for IC4651 No. 9122. The fitted orbit corresponds to a period of 747 days. A stellar jitter of  $15 \text{ m s}^{-1}$  has been added to the error bars.

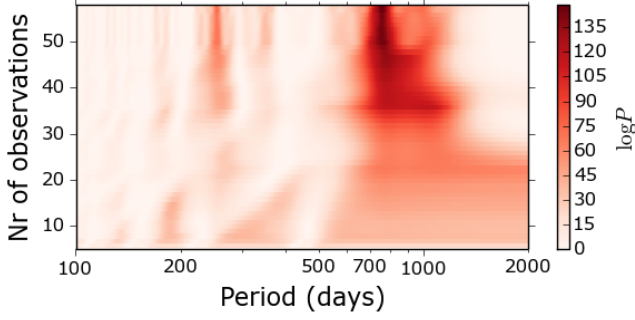


**Figure 4.** GLS of V magnitude for IC4651No9122 using ASAS data. The dashed line indicates the FAP at 1% level.

the planet candidate. In Fig. 4 we can see that there are two signals just above the FAP = 1% line with periods of 296 and 2273 days. Since we cannot constrain well the rotational period of the star it is difficult to claim whether any of these signals in the photometry correspond to such period.

### 3.4. Stellar activity and line profile analysis

To infer about activity modulations that could interfere with RV we measured a simultaneous activity index based on the H $\alpha$  line. We choose not to use the Ca II H&K and Na I D<sub>1</sub> & D<sub>2</sub> lines due to very low S/N ( $< 3$ ) on the former and line contamination on the latter. The H $\alpha$  index was calculated as in

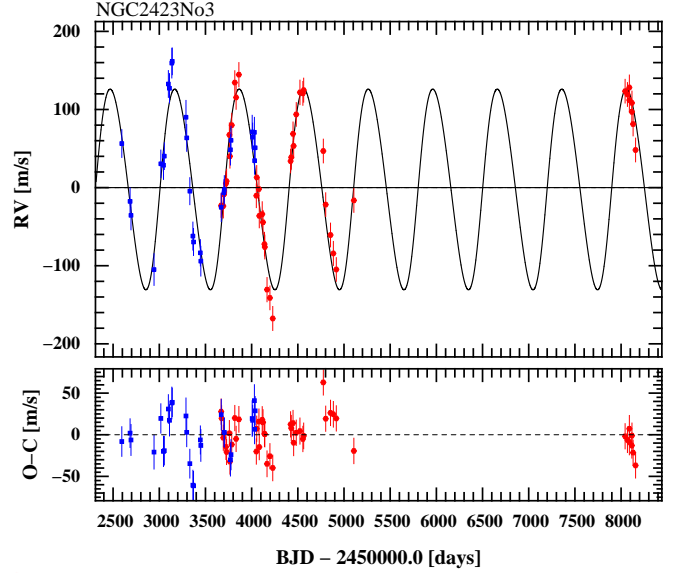


**Figure 5.** Stacked periodogram for RV measurements in IC4651 No. 9122.

Gomes da Silva et al. (2011) and implemented using ACTIN<sup>4</sup>. The flux in the  $H\alpha$  line was measured at 6562.81 Å with a 1.6 Å bandpass. Two reference lines were measured using two 10.75 Å and 8.75 Å windows centred at 6550.87 Å and 6580.31 Å, respectively. We then divided the flux in the  $H\alpha$  line by the two reference lines to obtain the index. The errors on the fluxes were computed using photon noise,  $\sqrt{N}$ , with  $N$  being the total flux in the band, and the error on the index was computed via propagation of errors. To further investigate RV variations caused by stellar atmospheric phenomena, we also used the full-width-at-half-maximum (FWHM) of the cross-correlation function (CCF). These values and their errors are provided by the HARPS pipeline. Moreover, we also analysed the Bisector Inverse Slope (BIS, Queloz et al. 2001) of the CCF but we note that this diagnostic of line asymmetry loses sensitivity for very low stellar rotations such as the case we are studying here (e.g. Saar et al. 1998; Santos et al. 2003; Queloz et al. 2009; Santos et al. 2014).

A periodogram analysis was carried out for FWHM, BIS and  $H\alpha$  index to infer about possible impact of stellar activity or long-term oscillations on the observed RV (see Fig. 2). The  $H\alpha$  periodogram shows two peaks at 952 and 714 days but they are not statistically significant. Interestingly, there are also peaks at 952 days in the BIS and RV periodograms although they are not significant. On the other hand, the FWHM periodogram also shows a long-period signal at  $\sim 714$  days, above the FAP = 1% line. This signal is probably due to rotational modulation of active regions in the atmosphere (see e.g. Larson et al. 1993; Lambert 1987) since it matches the same period as  $H\alpha$  index. As explained in Sect. 3.1, the determination of the rotation period is a very degenerate problem (depending on the assumption for macroturbulent velocity) and it is difficult then to determine whether that signal in the FWHM periodogram is caused by spot modulation or not. Nevertheless, we note that it has been reported that stars can have long-term activity cycles on top of the typical rotation-modulated cycle (Dumusque et al. 2011; Lovis et al. 2011). The fact that the period of the FWHM lies so close to the period of the planet candidate (which in turn matches one of the peaks of the  $H\alpha$  index periodogram) casts doubts on the planetary hypothesis and will be further discussed in Sect. 6. We note that during the referee process of this paper the work by Leão et al. (2018) announced the presence of a planet around IC4651 No. 9122 using the same RV data presented here. However, in that work the variation of the FWHM

<sup>4</sup> ACTIN is a FITS extractor and Activity Index calculator tool implemented in Python and available at <https://github.com/gomesdasilva/ACTIN>.



**Figure 6.** Radial velocity curve as a function of time for NGC2423 No. 3 using the complete dataset. The fitted orbit corresponds to a period of 698 days. A stellar jitter of 15 m s<sup>-1</sup> has been added to the error bars.

of the CCF was not explored which made the authors conclude that the signal should be caused by a planet. Our analysis casts doubts on the planetary nature of the periodic signal and shows the importance of analysing all possible stellar indicators.

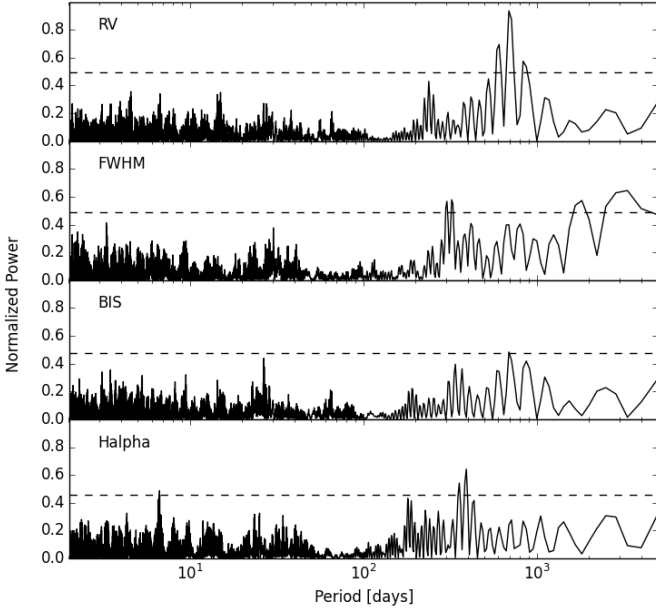
Stellar activity and pulsations are also expected to cause deformations in the line profile of spectral lines. Therefore, we also explored any possible correlation between the RV and several line-profile indicators such as FWHM, Bisector Inverse Slope (BIS, Queloz et al. 2001),  $V_{span}$  (Boisse et al. 2011), bi-Gauss (Nardetto et al. 2006) and  $V_{asy}$ , BIS+ and BIS- (Figueira et al. 2013) by using the python code *LineProf.py*<sup>5</sup> developed by Figueira et al. (2013). We did not find any strong correlation for any of the indicators, with a maximum value of the Pearson correlation coefficient of 0.3 for the relation of  $V_{span}$  with RV.

To further evaluate the significance of the periodic signals for RV we calculated the Stacked Bayesian general Lomb-Scargle (BGLS) periodogram which is shown in Fig. 5. This tool, developed by Mortier & Collier Cameron (2017), aims at tracking the S/N of the detection of a signal over time. If an orbiting body, such as a planet, is the cause of the observed periodic signal, we expect that the signal gets more significant as we add more observations. On the other hand, stellar activity signals are unstable and incoherent and thus the Stacked BGLS will show that the probability of having a signal at a given period varies (decreasing or increasing) when adding more data. In Fig. 5 we can see that the peak at  $\sim 700$  days gets more significant by adding more data and stabilises after  $\sim 50$  observations with some small fluctuations. This would be the expected behaviour for the signature of a planetary companion.

#### 4. Is the RV variation of NGC2423 No. 3 due to the presence of a planet?

In Paper I, showing the first results of this survey, a planet in the cluster NGC2423 was announced after collecting 46 RV points

<sup>5</sup> <https://bitbucket.org/pedrofigueira/line-profile-indicators>



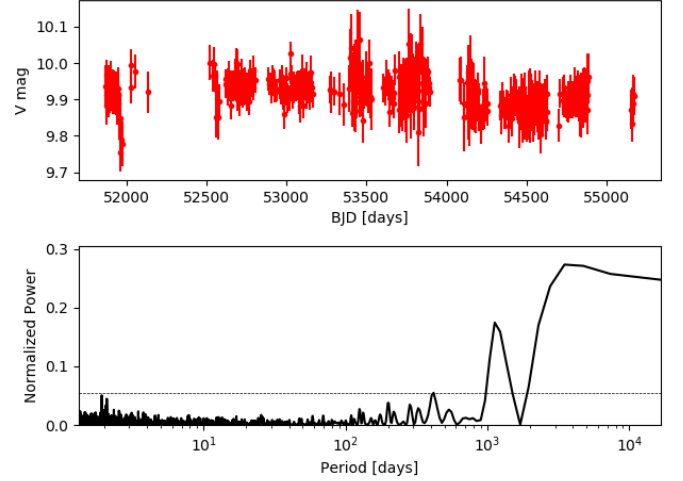
**Figure 7.** Generalized Lomb-Scargle periodograms of  $RV$ , FWHM, BIS and  $H\alpha$  index for NGC2423 No. 3. The dashed line indicates the FAP at 1% level.

(28 with CORALIE and 18 with HARPS) during 1529 days. The reported planet had a mass of  $10.6 M_J$  orbiting around NGC2423 No. 3 with a period of 714.3 days in an orbit of eccentricity 0.21 and semimajor axis 2.1 AU. The parent star has a mass of  $2.26 \pm 0.07 M_\odot$  and seems to be ascending the RGB and it is probably the most evolved star among the surveyed ones within this cluster. Most of the stars in this cluster exhibit a  $RV$  scatter below  $30 \text{ m s}^{-1}$  while NGC2423 No. 3 clearly contrasts with a dispersion of  $72 \text{ m s}^{-1}$ . We refer the reader to Paper I for the HR diagram and  $RV$  jitter histogram for this cluster. The parameters of NGC2423 No. 3 are detailed in Table 1. For this star we can also estimate the rotational period from the projected rotational velocity. If we consider a macroturbulence velocity of  $2.15 \text{ m s}^{-1}$  (with the formula by Gray 2005) we obtain  $v \sin i = 3.84 \text{ km s}^{-1}$ . Given that the radius of this star is  $13 R_\odot$  we would obtain a maximum period of 172 days. On the other hand, if we consider the empirical relation of Hekker & Meléndez (2007), the macroturbulence velocity would be  $5 \text{ km s}^{-1}$  which leads to a very small  $v \sin i$  ( $0.1 \text{ km s}^{-1}$ ) which in turn corresponds to a maximum rotational period of 6627 days.

#### 4.1. The radial velocity dataset

In Fig. 6 we show the  $RV$  data for NGC2423 No. 3 spanning more than 15 years of observations (73 observations, 45 of them with HARPS) and the orbital solution obtained with *yorbit* for a planet with a period of 698.6 days in a circular orbit with  $a = 2.02 \text{ AU}$ . The typical photon noise dominated errors in  $RV$  of the data is  $\sim 3 \text{ m s}^{-1}$  and  $\sim 10 \text{ m s}^{-1}$  for HARPS and CORALIE data, respectively. We note that we have also added here an error of  $15 \text{ m s}^{-1}$  before fitting the data<sup>6</sup>. By using the updated stellar mass ( $2.26 \pm 0.07 M_\odot$ ) from Delgado Mena et al. (2016) we derive a planetary mass  $m_2 \sin i = 9.6 M_J$  (see Table 2). The period

<sup>6</sup> In Paper I an error of  $20 \text{ m s}^{-1}$  was added, which was shown later to be overestimated, thus we decided to use here  $15 \text{ m s}^{-1}$  as done for the previous star.



**Figure 8.** GLS of  $V$  magnitude for NGC2423No3 using ASAS data. The dashed line indicates the FAP at 1% level.

and semi-major axis are very similar to previous values but the eccentricity is significantly lower (the former value was 0.21). The reduced  $\chi^2$  of the fit is 2.13 and the dispersion of the residuals is  $22.4 \text{ m s}^{-1}$ , a value similar to the typical error bars in  $RV$  considering the stellar jitter. The residuals shown in Fig. 6 seem to have a periodic behaviour. A GLS periodogram of the residuals shows two significant peaks around 400 days. Interestingly, this is close to the period of 417 days seen in the photometry (see below). However, we cannot properly fit such residuals with a single keplerian orbit or a sinusoidal function. We note here that by doing a two Keplerian fit, for the highest amplitude signal we would obtain similar orbital parameters as for the single Keplerian while the secondary signal would present a very bad fit. This probably reflects the fact that the residuals we see in the single Keplerian are of stellar origin, caused by different processes with different periods that our data cadence cannot resolve and thus cannot be fitted.

#### 4.2. Photometry

We collected 497 photometric measurements from the ASAS-3 catalogue (with grade A or B) taken between 2000 and 2009 with average errors of 0.04 mag. The lightcurve shows a  $\sim 0.15$  peak-to-peak variability in  $V$  magnitude. After applying a GLS we can see in Fig. 8 a clear peak above the FAP = 1% line with a period of 1124 days. Moreover, there is a second peak just above the FAP line at 417 days. Neither of these periods is close to the  $RV$  period but they could be related to the rotation of the star.

#### 4.3. Stellar activity and line profile analysis

In order to find any sign of stellar origin in the  $RV$  variation for this star we performed a similar analysis as for IC4651 No. 9122. In Figs. A.3 and 7 we show the time series and periodograms of the HARPS data only. The quality and resolution of CORALIE data does not allow to investigate the BIS for this star. For this star we could measure the  $\text{Na I D}_1$  &  $\text{D}_2$  lines but they do not show any periodic variation. The  $RV$  periodogram (Fig. 7) shows a strong peak at  $\sim 689.8$  days, similar to the period found with the Keplerian fit. The FWHM periodogram presents a signifi-

**Table 2.** Orbital and physical parameters for the planet candidates.

		IC4651 No. 9122	NGC2423 No. 3	NGC4349 No. 127
$\gamma$	[km s <sup>-1</sup> ]	-30.2868 ± 0.0107	18.3014 ± 0.0078	-11.4322 ± 0.0119
P	[days]	747.22 ± 2.95	698.61 ± 2.72	671.94 ± 5.32
K	[m s <sup>-1</sup> ]	100.765 ± 8.427	128.451 ± 5.851	229.454 ± 9.932
$e$		0.150 ± 0.068	0.088 ± 0.041	0.046 ± 0.033
$\omega$	[deg]	-141.9 ± 16.1	-102.1 ± 23.2	168.3 ± 54.5
Tp	[BJD-2 400 000]	55575.85 ± 36.31	55083.27 ± 43.16	54306.62 ± 100.64
$m_2 \sin i$	[M <sub>J</sub> ]	7.202	9.621	24.097
$a$	[AU]	2.05	2.02	2.35
$N_{meas}$		57	73	46
Span	[days]	4592	5563	1587
$\Delta v$ (HARPS-Coralie)	[km s <sup>-1</sup> ]	—	0.007	—
$N\sigma$	[m s <sup>-1</sup> ]	17.00	22.40	31.79

cant peak at  $\sim 322$  days with an alias at  $\sim 303$  days caused by the long gap in time between the two sets of observations<sup>7</sup>. The same happens to the H $\alpha$  periodogram, with a stronger peak at  $\sim 392$  days and an alias at  $\sim 357$  days also above the FAP level. This signal might be related to the period of 417 days observed in the photometry and be produced by rotational modulation of active regions. Finally, the BIS measurements show a peak just above the FAP line with exactly the same period as the  $RV$  variability, 689.8 days.

We also studied a possible relation between the  $RV$  and line-profile indicators. In Fig. A.4 we can see how the  $RV$  of NGC2423 No. 3 shows a strong correlation with biGauss and BIS (the Pearson's correlation coefficient is depicted in each plot). This fact warns us about the possibility that the signals we are observing are related to inhomogeneities in the stellar surface and not with an orbiting body. However, we note that the 18  $RV$  points initially presented in Paper I did not show any correlation with BIS and in some of the observing time windows the BIS is practically flat and thus not correlated with  $RV$  (see Fig. A.3). It is interesting however, that the sign of the slope for BIS vs  $RV$  is opposite to what is found in many cases in the literature (e.g. Queloz et al. 2001; Boisse et al. 2009; Figueira et al. 2013), as expected for  $RV$  variations due to stellar spots. This case resembles that of the slow rotator HD 41004A (Santos et al. 2002) where BIS shows a strong positive correlation with  $RV$ . In that system the best explanation for the  $RV$  variation found and the positive correlation with BIS is given by the presence of a brown dwarf around the secondary component HD 41004B (probably an M2 dwarf), which perturbs the contribution to the global CCF dominated by the primary. However, in our case, there are not reported binary companions around the planet candidate host but we cannot discard the possibility that there is a blended binary. By inspecting the SIMBAD database, we find that a foreground star from the cluster NGC2422 (PMS 419) with  $V=18.5$  is at 9.18 arcsec distance from our star. That cluster is closer and younger than NGC2423 (490 pc vs 766 pc as in WEBDA database). By using the simulations done by Cunha et al. (2013) we can estimate a contamination of  $\sim 1$  m s<sup>-1</sup> (considering that our target is a K5 star and the difference in apparent magnitude is 8). Thus, we can reject the hypothesis that such foreground star is affecting our  $RV$  measurements. On the other hand we do not

find a correlation of  $RV$  with FWHM and contrast as also happens in the case of HD 41004. Moreover, Santerne et al. (2015) simulated different blending scenarios (planet orbiting around a secondary unseen companion) and showed that when the secondary rotates faster than the primary ( $FWHM_1 < FWHM_2$ ) a positive correlation between BIS and  $RV$  would be observed, as in our case. However, a positive correlation would also be observed for BiGauss and  $V_{span}$  meanwhile here we do find a negative correlation for those two line profile indicators. Another case for a positive correlation between BIS and  $RV$  was found for the active K giant EK Eri (Dall et al. 2005) which later was reported to have a strong cool spot by Aurière et al. (2011) using the Zeeman-Doppler imaging method.

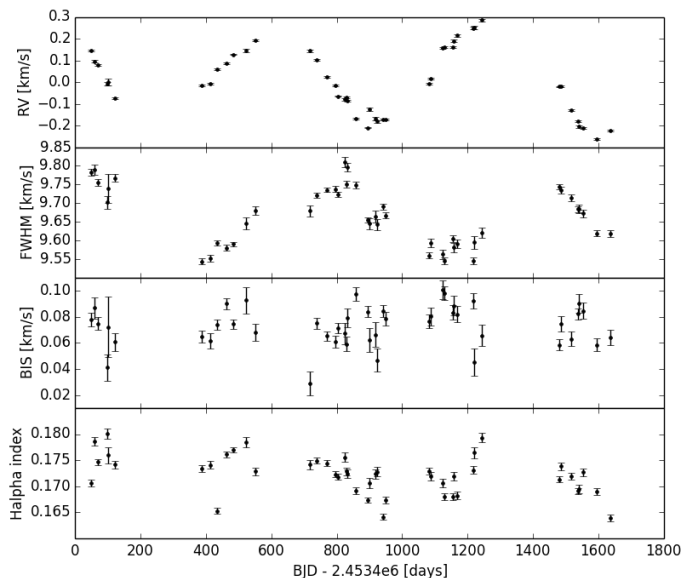
The real value of  $v \sin i$  is uncertain, as in the case of IC4651 No. 9122. Thus, it is difficult to relate the periodicity of the FWHM with rotational modulation of active regions but we cannot discard that possibility. Indeed, the different peaks observed in the periodograms of FWHM, H $\alpha$  and photometry (322, 392 and 417 days), on one hand, and of BIS (689.8 days) on the other hand, might indicate that the variations of the former are due to rotational modulation whereas the period of the latter is related to stellar pulsations which in turn can also be producing the periodic variation in  $RV$ . Indeed, the fact of having a signal with eccentricity close to zero cannot rule out the possibility that pulsations are causing the  $RV$  variation (Frink et al. 2002). For this star we find large peak-to-peak  $RV$  amplitudes of  $\sim 270$  m s<sup>-1</sup> meanwhile BIS variations are much smaller (40 m s<sup>-1</sup>), which could be a signature of oscillations (Hatzes 1996).

In an attempt to further understand the different signal observed in this star we show the stacked BGLS periodogram of  $RV$  in Fig. A.5. The probability of the signal steadily increases by adding new observations, thus the signal is stable.

## 5. A stellar signal disguising as a brown dwarf around NGC4349 No. 127

In Paper I, a 19.8 M<sub>J</sub> substellar companion around NGC4349 No. 127 was presented with a period of 677.8 days,  $e = 0.19$ ,  $K = 188$  m s<sup>-1</sup> and  $a = 2.38$  AU. As for the previous case, this star is also the most evolved in the cluster (see also Paper I and Delgado Mena et al. 2016, for HR diagrams and  $RV$  scatter histogram) and it clearly stands out with a  $RV$  scatter of 70 m s<sup>-1</sup> (the average  $RV$  jitter is 20 m s<sup>-1</sup>). The estimated  $v \sin i$  for this star is 6.1 km s<sup>-1</sup> (assuming a macrotubulence velocity of

<sup>7</sup> We note that we have corrected the FWHM values since they tend to slightly increase with time due to the degradation of the spectrograph focus but still an increasing drift can be observed.



**Figure 9.** Time-series of  $RV$ , FWHM, BIS and  $H\alpha$  index for NGC4349 No. 127.

$2.1 \text{ km s}^{-1}$  derived from Valenti & Fischer 2005) and the radius is  $37 R_{\odot}$  which leads to a maximum rotational period of  $\sim 307$  days. However, as shown in previous sections, if we consider the macroturbulence given by Hekker & Meléndez (2007) relation ( $4.8 \text{ km s}^{-1}$ ) we would obtain  $v \sin i = 4.8 \text{ km s}^{-1}$  and a maximum rotational period of  $\sim 390$  days. The stellar parameters for this star can also be seen in Table 1.

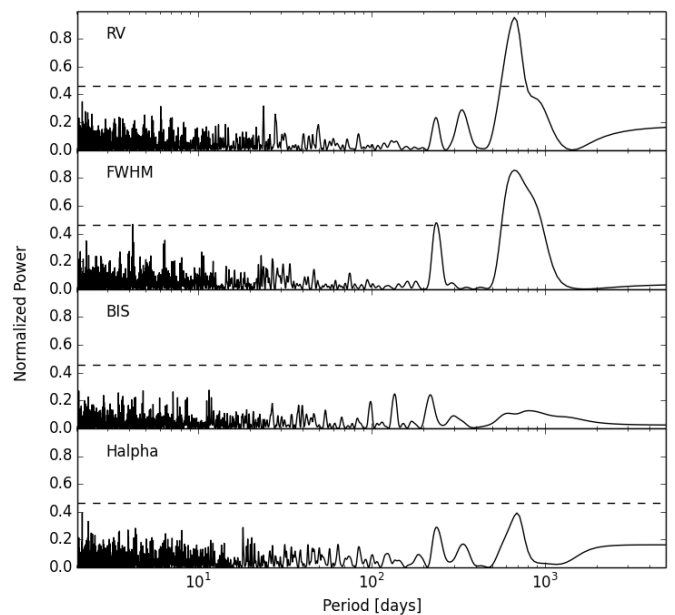
### 5.1. Radial velocity dataset

In Fig. 9 we show the time series for  $RV$ , FWHM, BIS and  $H\alpha$  for NGC4349 No. 127. For this star the  $\text{Na I D}_1$  &  $\text{D}_2$  lines are contaminated. This plot includes 46 measurements collected during 1587 days while in Paper I, 20  $RV$  measurements along 784 days were presented. Figure 10 shows the  $RV$ , FWHM, BIS and  $H\alpha$  periodograms for NGC4349 No. 127. The signal attributed to a brown dwarf candidate in Paper I has a significant peak at  $\sim 666$  days in the  $RV$  periodogram which is similar to the period that we also find if we fit a Keplerian orbit to the data (672 days, Fig. 11). This value is very close to the 678 days period derived in Paper I. We obtain a planetary mass  $m_2 \sin i = 24.1 M_J$  and the eccentricity value is 0.05, lower than previously reported value in Paper I.

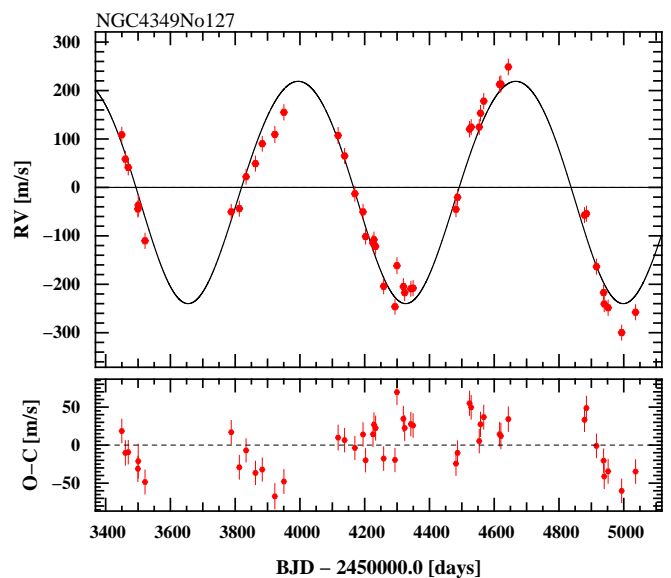
### 5.2. Stellar activity, photometry and line profile analysis

In Figure 10 we can see how the FWHM periodogram also shows a very strong peak at  $\sim 666$  days (middle panel), indicating that there are variations in the star’s atmosphere with the exact same period as the  $RV$ . Moreover, the  $H\alpha$  periodogram also shows a signal at a similar period, with  $P \sim 689$  days (however just below the  $FAP = 1\%$  level). We can also see a significant signal at 235 days for the FWHM which has a counterpart in the  $H\alpha$  periodogram but with a lower power.

The time-series of  $RV$ , FWHM, and  $H\alpha$  show clearly the signals detected in the respective periodograms (see Fig. 10). We note that the amplitude of the FWHM variation is significant (peak-to-peak variation of  $\sim 250 \text{ m s}^{-1}$ ). In the figure, we can observe that FWHM and  $RV$  are not in phase. However, stellar



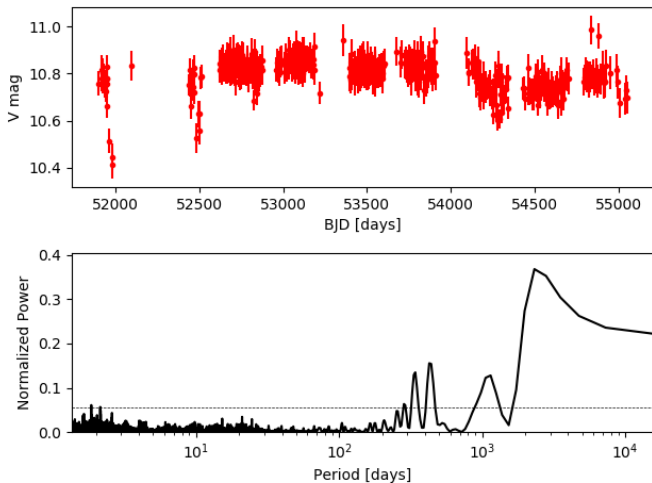
**Figure 10.** Generalized Lomb-Scargle periodograms of  $RV$ , FWHM, BIS and  $H\alpha$  index for NGC4349 No. 127. The dashed line indicates the FAP at 1% level.



**Figure 11.** Radial velocity data as a function of time for NGC4349 No. 127. The fitted orbit corresponds to a period of 672 days. A stellar jitter of  $15 \text{ m s}^{-1}$  has been added to the error bars.

rotationally modulated active regions are known to induce  $RV$  signals not in phase with activity proxies (e.g. Queloz et al. 2001; Santos et al. 2014). In their detection paper, Lovis & Mayor (2007) used BIS and the  $\text{Ca II H\&K}$  lines as activity diagnostic tools to validate the candidate brown dwarf. However, BIS loses sensitivity for slow rotators and the stellar flux in the  $\text{Ca II H\&K}$  spectral region for red stars is very low, delivering a low S/N for this activity index. This could explain why the authors missed the strong signals coming from the star’s atmosphere with the same period as the observed  $RV$  variations.





**Figure 12.** GLS of V magnitude for NGC4349No127 using ASAS data. The dashed line indicates the FAP at 1% level.

For this star we found two independent sets of photometric measurements from the ASAS-3 catalogue. In one of the sets we do not find any significant peak in the GLS applied to the light curve. The second set is shown in Fig. 12. We find three significant signals in the periodogram with periods 428, 342 and 1149 days. The first two periods might be related to the rotational period of the star which we have estimated to be in between 300 and 400 days. However, there is no clear signal around the 666 days period found in the *RV* although we note that the first estimation of the maximum rotational period of  $\sim 307$  days and the signal in photometry at 342 days are close to the first harmonic ( $P/2$ ) of the *RV* period for NGC4349No127.

On the other hand, we have compared the *RV* with the line-profile indicators as done in previous section and we do not find significant correlations among the different indicators. Moreover, we show in Fig. A.6 the stacked BGLS for *RV*. The signal looks stable and increase its significance by adding more observations which would be the expected behaviour for an orbiting body but such a trend is also expected if the stellar phenomena responsible for the *RV* signal is stable over the time interval of our measurements.

## 6. Discussion

As we have seen in previous sections there can be different causes for the *RV* variability detected in our targets apart from the presence of planets. Long-period *RV* variations with hundreds of days have been known to exist in several giant stars (Walker et al. 1989; Larson et al. 1993; Hatzes & Cochran 1993) with *RV* amplitudes in the order of  $\sim 50$ – $400$   $\text{m s}^{-1}$  which were attributed to rotationally modulated active regions (Larson et al. 1993; Lambert 1987) or radial and non-radial pulsations (Hatzes & Cochran 1999), depending on the period of the signal. However, it is not clear yet whether pulsations (and of which kind) can create such high amplitude *RV* variations as the ones observed. Indeed, stochastically-driven p-mode oscillations (also known as solar-type pulsations) are observed in giants from the lower RGB all the way to the AGB evolutionary phase (Mosser et al. 2013), where the most evolved of these are known as semi-regular variables. The properties of these pulsations follow specific scaling relations that can be used to predict their period and amplitude. Considering the stellar parameters

estimated for the three stars under study, the expected periods would be less than a few days (the dynamical timescale for, e.g., NGC 2423 No. 3 is  $\sim 20.8$  hr). This, therefore, rules out the possibility that the variability seen in these stars is due to pulsations excited by a stochastic mechanism. On the other hand, if, by absurd, we were to assume that these type of pulsations justify the signal observed, and extrapolated from Eq. (9) of Mosser et al. (2013), we would find an rms variability amplitude of 0.78 mag for modes with periods of about 698 days (the case of NGC2423 No. 3), hence a peak visible magnitude change of 1.1, which is well above what we see here. The use of the scaling relation in Kjeldsen & Bedding (1995) would then point to an rms velocity amplitude of the order of  $10 \text{ km s}^{-1}$ , thus several orders of magnitude above what we see here.

Two other type of pulsations are observed in giants, that cannot be explained by a stochastically-driven mechanism. The first is the Mira variability, which shows much larger amplitudes than the oscillations discussed above. These are seen only at luminosities much higher than our target stars and, thus, can also be ruled out. The second - whose pulsation origin is still under debate (e.g. Trabucchi et al. 2017; Takayama et al. 2015) - has periods of several hundred days, significantly longer than the fundamental radial mode for the stars where they are observed. They appear as a particular sequence in the period-luminosity relation (sequence D) that is yet to be fully understood and are found both during the RGB and the AGB at luminosities  $\log(L/L_{\odot}) \gtrsim 2.5$  dex. Recently they have been proposed to be manifestations of oscillatory convective modes (Saio et al. 2015). While in current theoretical models these appear only in giants with  $\log(L/L_{\odot}) \gtrsim 3$  dex (for a  $2 M_{\odot}$  star), for models with reasonable convection parameters, the exact minimum luminosity depends on the mass and the mixing length adopted in the treatment of convection. NGC4349 No. 127 has  $\log(L/L_{\odot}) \sim 2.7$  and could be a candidate for showing this type of pulsations, but we note that the period of 672 days is larger than expected for a typical sequence D star, unless the luminosity is significantly underestimated. The other two stars under study have much lower luminosity, so the presence of this type of pulsations seems to be even less likely, given the theoretical models. Also, their periods would be too long to fit into the sequence D described before. Moreover, the models by Saio et al. (2015) predict large surface temperature variations that are however not observed in our stars<sup>8</sup>.

In summary, although pulsations are a possibility to explain the long-period variability seen in the *RV*, after considering the stellar parameters of the stars studied in this paper and the period and amplitude of the variability, we find that the type of pulsations discussed in the literature cannot comfortably explain our observations, with the possible exception of NGC4349 No. 127. We thus conclude, that either these are not produced by pulsations or they hint at a type of pulsation variability yet to be explained.

Alternatively, stellar activity and/or magnetic fields might be the cause for the *RV* variability observed in giant stars. Among the three cases exposed in this work, probably, the easiest to explain would be NGC 4349 No. 127 since we find a perfect match between the periods of the *RV*, the FWHM of the CCF and the  $H_{\alpha}$  indicator. The variability we observe in these three indicators is thus presumably caused by rotational modulation of active regions. However, we are still far from fully understanding how the magnetic fields behave in giant stars and what are the

<sup>8</sup> We have derived  $T_{\text{eff}}$  values for each individual spectrum and the maximum variation we observe is of  $\sim 50$  K with no clear periodicity.

timescales of stellar activity signals since their low rotation rates prevent their study through classical Doppler imaging, a successful technique to characterize the distribution of spots. Starspots with lifetimes of years have been observed with this technique in some stars (Strassmeier 2009; Hussain 2002). For example, polar cap-like starspots appear to have lifetimes of over a decade in RS CVn binaries as well as in young MS stars (Strassmeier 2009). Nevertheless, a handful of giant stars with slow rotation have been studied in the last decade with the Zeeman Doppler imaging, which allows to obtain the strength and geometry of the surface magnetic field by using new high resolution spectropolarimeters such as Narval, ESPaDOnS or HARPSpol (Aurière et al. 2015). These authors found that most of the giant stars with detected magnetic fields (the longitudinal component) are undergoing the first dredge-up (i.e. they are close to the RGB base) and few of them are burning He in the core (red clump). Moreover, stars with rotational periods shorter than 200 days seem to have magnetic fields driven by a dynamo as in our Sun. Those stars present a good correlation between the strength of the magnetic field, the chromospheric activity ( $S$ -index) and the rotational period. On the other hand, stars such as EK Eri (Aurière et al. 2011) or  $\beta$  Ceti (Tsvetkova et al. 2013) with  $P_{rot}$  of 308.8 and 215 days, respectively, probably represent the prototype of descendants of magnetic Ap-Bp stars, hosting fossil magnetic fields. They have strong magnetic fields ( $|B_{lmax}| > 10$  G) with lower than expected chromospheric activity. Finally, a last group of stars present very weak (sub Gauss level) but detectable magnetic fields with higher than expected chromospheric activity. Among this last group there are several interesting cases such as Pollux, Aldebaran and Arcturus.

Pollux ( $\beta$  Gem) shows  $RV$  variations with a period of 590 days spanning 25 years (Hatzes et al. 2006) attributed to the presence of a planet. The star does not show any variability in the chromospheric  $S$ -index or in the bisector velocity span (Hatzes et al. 2006). The photometric period does not match either that of the  $RV$  variations. However,  $RV$  values are correlated with the longitudinal magnetic field (Aurière et al. 2009, 2014), casting doubts on the planetary nature of the signals. Aurière et al. (2015) also reported a weak magnetic field in Aldebaran ( $\alpha$  Tau) but they could not determine the period of its variability. An extensive study of  $RV$  in this star spanning 20 years by Hatzes et al. (2015) reports the presence of a giant planet with a period of  $\sim 630$  days and modulation of active regions with a period of  $\sim 520$  days. Finally, a weak magnetic field was also detected in Arcturus ( $\alpha$  Boo) by Sennhauser & Berdyugina (2011) with a possible period of 208 days (only 3 measurements were made). The detection was confirmed by Aurière et al. (2015) but no period is provided. Therefore, it remains unclear if the high amplitude  $RV$  variations ( $500 \text{ m s}^{-1}$ ) with  $P = 231$  days detected by (Hatzes & Cochran 1993) could be caused by the magnetic field. These three cases reflect very well the problematic of establishing the nature of different signals even when a large amount of data are collected during many years.

With the current data for the three stars studied here it is difficult to assure whether stellar chromospheric activity is affecting the  $RV$  measurements. When adding all the spectra for a given star to increase the S/N in the blue region we cannot appreciate any sign of emission in the cores of Ca II H&K lines. The semi-amplitude of the BIS has values of  $\sim 20 \text{ m s}^{-1}$  for IC4651 No. 9122 and NGC2423 No. 3 and  $\sim 30 \text{ m s}^{-1}$  for NGC4349 No. 127. Considering a maximum  $v \sin i$  of  $4 \text{ km s}^{-1}$  for the former and  $6 \text{ km s}^{-1}$  for the latter, the presence of spots with a filling factor higher than 5% would be needed to explain the BIS semi-amplitudes (see Fig. 4 of Santos et al. 2003). However, for val-

ues of  $v \sin i$  of  $2 \text{ km s}^{-1}$  or lower the presence of spots would not produce BIS variations. A spot filling factor of 5%-10% would produce a variability in  $V$  of  $\sim 0.05$ - $0.11$  mag which is below the peak-to-peak variability we observe for our stars (with approximate values of 0.3, 0.15 and 0.2 mag for IC4651 No. 9122, NGC2423 No. 3 and NGC4349 No. 127, respectively). Therefore, the photometric variability we observe in our stars might be caused by spots (which in turn would produce the observed BIS variations) but the photometric periods do not match the period of the  $RV$  neither the period of the BIS for any of the stars.

We find striking the fact that the periods in  $RV$  found for the three stars are all close to 700 days. We also find hints that the amplitude and the phase of the signal is changed along the time for the three cases. Curiously, in a recent paper by Hatzes et al. (2018) they also report a  $RV$  signal in  $\gamma$  Draconis (a K giant with  $T_{\text{eff}} = 3990 \text{ K}$ ,  $\log g = 1.67$ ,  $2.14 M_{\odot}$  and  $L = 510 L_{\odot}$ ) with a semi-amplitude of  $148 \text{ m s}^{-1}$  and a period of 702 days that changes in phase and amplitude. It is interesting, though, that they do not find any correlation of  $RV$  with the Ca II index or with the BIS. These authors suggest that the variability observed in  $\gamma$  Draconis might be caused by dipole oscillatory convection modes (Saio et al. 2015) which could be also the explanation for the  $RV$  variability observed in NGC4349 No. 127 as discussed above. There are other cases in the literature with long-period  $RV$  variations in M giants as well. For example, Lee et al. (2017) reported that the M giant HD 36384 shows photometric variations with a period of 570 days close to the period of 535 days in the  $RV$ . Furthermore, the M giant  $\mu$  UMa shows  $RV$  variations with a period of 471 days with bisector velocity curvature variations of 463 days (Lee et al. 2016). Moreover, the EWs of H $\alpha$  and H $\beta$  lines show periodic variations with 473 days for this star. Recently, Bang et al. (2018) found  $RV$  variations with a period of 719 days in the M giant HD 18438 which also shows photometric and H $\alpha$  variability with similar period. All these examples in the literature together with the three stars presented here warns us about the difficulty of confirming the planetary nature of periodic  $RV$  variations in red giants, specially the most massive ones.

In the last years several works have discussed the planet occurrence as a function of stellar mass finding different results (e.g. Omiya et al. 2009; Johnson et al. 2010; Reffert et al. 2015; Jones et al. 2017). When Paper I was published, only 6 substellar companions<sup>9</sup> were known around stars more massive than  $\sim 1.78 M_{\odot}$  and they presented much higher masses on average than for lower mass stars (see Fig. 11 in Paper I), leading to the hypothesis that the frequency of massive planets is higher around more massive stars. Certainly, the inclusion of the brown dwarf around the massive star NGC4349 No. 127 contributed to such correlation. Later studies by the Korean–Japanese planet search program confirmed such result (Omiya et al. 2012). The works by Reffert et al. (2015); Jones et al. (2016) also suggest that the planet occurrence rate increases with stellar mass but peaks around  $2 M_{\odot}$  and sharply decreases for stars more massive than  $2.5$ - $3 M_{\odot}$  in agreement with the theoretical model by Kennedy & Kenyon (2008). Indeed, the 12-years length Lick Observatory survey (Reffert et al. 2015) did not find any planet around stars more massive than  $2.7 M_{\odot}$  despite having a significant number of stars with masses up to  $5 M_{\odot}$ . Similarly, the 6-years length EXPRESS survey has not found any planet in stars more massive than  $2.5 M_{\odot}$  (Jones et al. 2016). Nowadays, in the

<sup>9</sup> Nowadays, this number goes up to 67 companions by using the data collected in The Extrasolar Planets Encyclopaedia

Extrasolar Planets Encyclopaedia we can find only 8 stars with masses above  $2.7M_{\odot}$  hosting planets or brown dwarfs<sup>10</sup>, being HD 13189 the only one more massive than NGC4349 No. 127, which should now be excluded. Moreover, the two bodies system around the  $2.8M_{\odot}$  star BD+20 2457 seems to be dynamically unstable (Horner et al. 2014), thus, the determination of planet occurrence rates with surveys covering the high stellar mass regime is advisable. As mentioned in the introduction, the measurement of masses for giant stars is a difficult task and several works discuss about the suitability of using isochrones compared with other methods such as asteroseismology which delivers much more precise masses. In order to do a proper statistical analysis evaluating planet occurrence rates we need to have a sample of homogeneously derived masses. Such study is out of the scope of the present paper but will be approached in a future work when the data collection for our survey in open clusters is completed.

## 7. Conclusions

In this work we present the *RV* data for three stars within the  $\sim 15$  years survey of radial velocities in a sample of more than 142 giant stars in 17 open clusters. Additional data was collected during the last year. The large span of the data is needed in order to find long period planets since the probability of discovering short period planets around such kind of stars diminishes as the stars evolve and expand their radii. The long-term monitoring of giant stars is also essential to distinguish *RV* variations caused by the rotational modulation of active regions because the rotational periods of evolved stars are much longer than for dwarfs.

We report the possible discovery of a  $7.2M_J$  planet in a 747 days orbit around the star No. 9122 in the intermediate-age open cluster IC4651. If confirmed, the planet would be at a distance of 2.05 AU and the eccentricity of the orbit would be 0.15. However, we find that the FWHM values of the CCF and the  $H\alpha$  index show a periodic behavior with a bit shorter period than the *RV* values. The period of the FWHM and  $H\alpha$  index variation get even closer to that of the *RV* when adding few more points taken after 2009. Nevertheless, we note that the significance of the periodic signals on those two indicators is low. We perform further tests on this system to understand whether the periodic *RV* signals are of stellar origin. Although there exists the possibility of having both rotational modulation and a planet with close periods we conclude that more data will be needed to validate or refute the presence of such planet.

This survey already reported two substellar companions in Paper I, NGC2423 No. 3b and NGC4349 No. 127b. These stars were further followed in the next two years after the publication of Paper I (for NGC2423 No. 3 additional observations were carried out during last semester) and here we present an updated analysis of the *RV* data. We find that the *RV* variability of NGC2423 No. 3 is strongly correlated with the BIS of the CCF when using the full dataset, a fact that by itself probably invalidates the planet presence although we cannot explain the reason for such correlation. However, we note that such correlation is not stable along the time and the BIS becomes flat in some of the observed time windows. On the other hand, for NGC4349 No. 127 we find that the FWHM of the CCF presents a large amplitude periodic variation with the same strong peak as

the *RV*. Moreover, the periodogram of the  $H\alpha$  index also shows a peak with a similar period very close to the 1% FAP level. Therefore, we believe that the presence of active regions in the atmosphere and not a substellar companion is responsible for the *RV* variability. For this star, the presence of a new kind of stellar pulsations might be a plausible explanation for the observed variability as well.

Long-term period signals have been observed in many K giants either produced by stellar activity or by oscillations. The amplitudes of these signals can also reach hundreds of  $\text{m s}^{-1}$ , mimicking the presence of giant planets. The three cases presented in this work are a clear example of the need to perform detailed analysis of *RV* modulations in giant stars before the real nature of a planetary signal can be assumed. Moreover, the analysis of the FWHM of the CCF has revealed as essential in order to discard the presence of a companion but we note that this indicator is not widely used in the validation of planets around evolved stars. Finally, more theoretical and observational study is needed to further understand the different nature of stellar oscillations in evolved stars.

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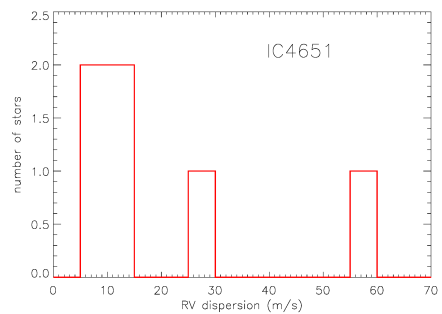
This research has made use of the The Extrasolar Planets Encyclopaedia, SIMBAD and WEBDA databases. This work has also made use of the IRAF facility.

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<sup>10</sup> We note that not all the planet hosts have an entry for the stellar mass in this database and the available values listed there come from different sources/methods, thus any conclusion made from this heterogeneous dataset must be taken with caution.

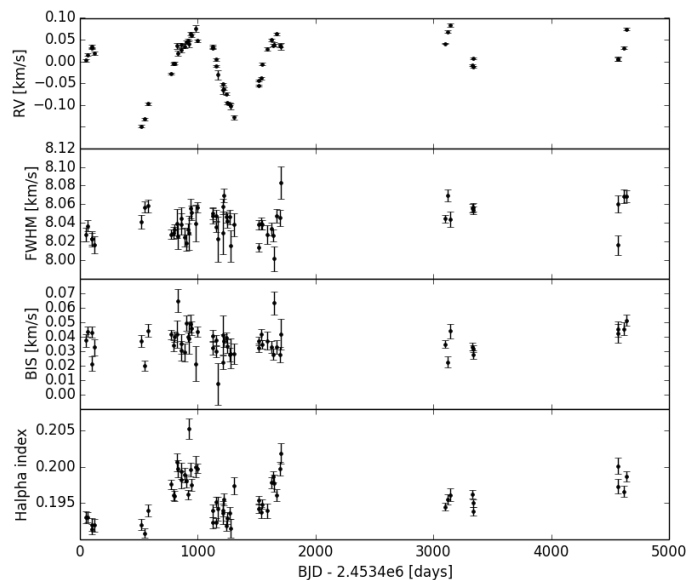
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**Figure A.1.** Standard deviation of *RV* for members in the cluster IC4651. Two stars with dispersions higher than  $100 \text{ m s}^{-1}$  (due to stellar or brown dwarf companions) are not shown here.

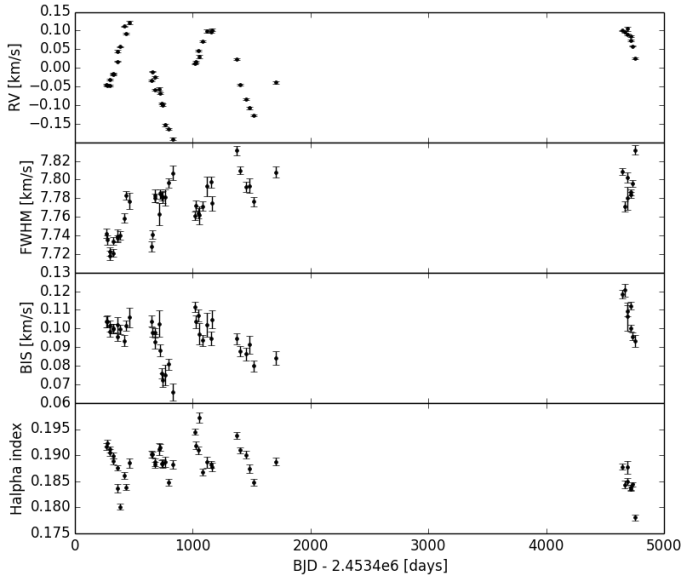
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## Appendix A: Additional figures

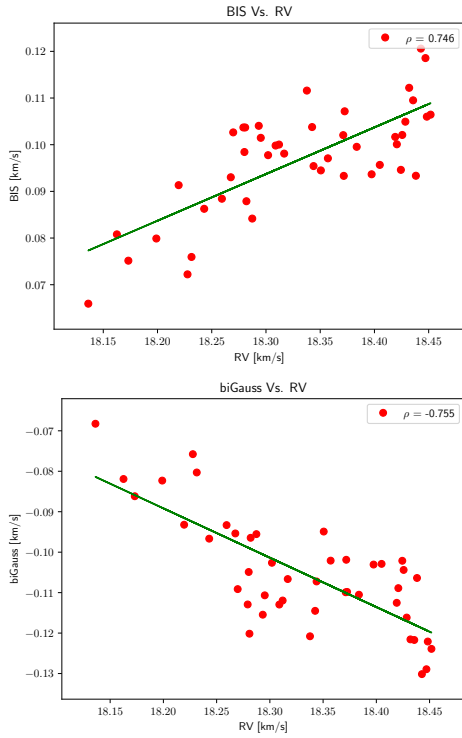


**Figure A.2.** Time series of *RV*, FWHM, BIS and  $H\alpha$  index for IC4651 No. 9122.

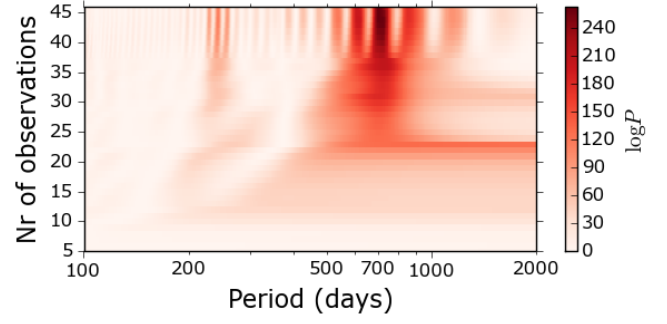
## Appendix B: *RV* measurements



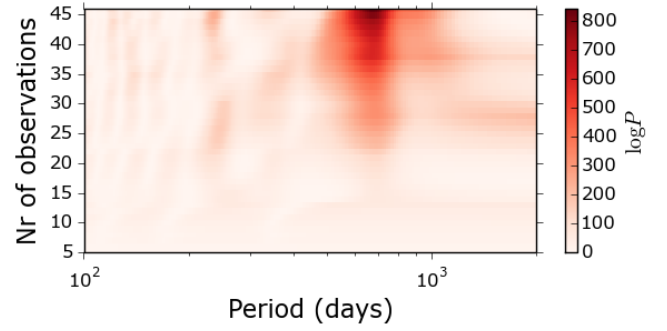
**Figure A.3.** Time series of  $RV$ , FWHM, BIS and  $H\alpha$  index for NGC2423 No. 3.



**Figure A.4.** Correlation between  $RV$  and the line profile indicators BIS and biGauss for NGC2423 No. 3.



**Figure A.5.** Stacked periodogram for  $RV$  measurements in NGC2423 No. 3.



**Figure A.6.** Stacked periodogram for  $RV$  measurements in NGC 4349 No. 127.

**Table B.1.** Radial velocity measurements for IC4651 No. 9122 obtained with HARPS and their instrumental error bars. All data are relative to the solar system barycenter.

JD-2 400 000	<i>RV</i> [km s <sup>-1</sup> ]	Uncertainty [km s <sup>-1</sup> ]
53449.887674	-30.25225	0.00317
53462.850455	-30.23935	0.00257
53499.742442	-30.22078	0.00363
53500.698475	-30.22416	0.00272
53521.693417	-30.23617	0.00397
53921.668152	-30.40372	0.00302
53950.621071	-30.38763	0.00256
53981.622808	-30.35229	0.00301
54171.886988	-30.28285	0.00231
54194.857714	-30.25975	0.00282
54200.849131	-30.25867	0.00267
54226.894524	-30.21881	0.00658
54232.731181	-30.23582	0.00561
54256.735797	-30.21726	0.00492
54261.807071	-30.22907	0.00488
54291.722147	-30.21941	0.00428
54299.690631	-30.20983	0.00379
54319.614198	-30.20572	0.00244
54323.579256	-30.21409	0.00786
54342.555162	-30.19083	0.00437
54349.595115	-30.19470	0.00322
54385.490351	-30.17916	0.00851
54393.518352	-30.20727	0.00239
54523.892742	-30.22074	0.00341
54529.878124	-30.22392	0.00284
54553.859622	-30.25017	0.00255
54556.860227	-30.26548	0.00240
54568.915197	-30.28588	0.01020
54615.900828	-30.32074	0.00974
54616.851591	-30.30711	0.00347
54621.885865	-30.31795	0.00317
54641.777674	-30.32915	0.00254
54646.614534	-30.35011	0.00332
54672.612166	-30.35315	0.00332
54679.582438	-30.35829	0.00721
54708.574369	-30.38375	0.00517
54913.858276	-30.31120	0.00213
54919.803228	-30.29862	0.00200
54936.812674	-30.29323	0.00286
54949.768030	-30.26058	0.00211
54990.708393	-30.22648	0.00427
55024.810993	-30.20521	0.00291
55038.738466	-30.21751	0.00282
55045.784236	-30.21564	0.00559
55072.564309	-30.19168	0.00319
55095.574817	-30.21826	0.00370
55105.500100	-30.21990	0.00749
56499.579896	-30.21348	0.00184
56521.628162	-30.18630	0.00278
56543.505489	-30.17094	0.00352
56733.852100	-30.26437	0.00198
56734.867908	-30.24730	0.00200
56735.863877	-30.26788	0.00173
57968.650782	-30.25014	0.00400
57969.606900	-30.24899	0.00450
58013.569803	-30.22394	0.00317
58041.498430	-30.18038	0.00278

**Table B.2.** Radial velocity measurements for NGC2423 No. 3 obtained with HARPS and their instrumental error bars taken after JD-2 454 130. Older data is presented in Paper I. All data are relative to the solar system barycenter.

JD-2 400 000	<i>RV</i> [km s <sup>-1</sup> ]	Uncertainty [km s <sup>-1</sup> ]
53669.846074	18.28083	0.00215
53674.796790	18.27917	0.00224
53692.862440	18.28013	0.00217
53699.842505	18.29520	0.00210
53721.855913	18.30887	0.00187
53728.752333	18.31202	0.00191
53758.656817	18.37122	0.00301
53764.700951	18.34386	0.00150
53784.639237	18.38365	0.00190
53817.563704	18.43820	0.00214
53831.568254	18.41915	0.00195
53861.565085	18.44831	0.00369
54050.824982	18.29334	0.00226
54054.864236	18.31679	0.00191
54078.820691	18.30194	0.00190
54082.782001	18.26757	0.00288
54114.702865	18.26979	0.00499
54122.715976	18.25936	0.00235
54138.717872	18.23135	0.00197
54142.652470	18.22766	0.00270
54166.590019	18.17312	0.00385
54197.574062	18.16256	0.00218
54228.532429	18.13619	0.00333
54421.769842	18.33770	0.00209
54428.783945	18.34252	0.00248
54447.806766	18.37255	0.00234
54454.813715	18.35709	0.00399
54482.750391	18.39740	0.00220
54522.713892	18.42563	0.00452
54554.593644	18.42438	0.00256
54563.523609	18.42850	0.00332
54775.842423	18.35059	0.00208
54801.736645	18.28197	0.00193
54853.716032	18.24304	0.00251
54884.669017	18.21949	0.00333
54916.569668	18.19892	0.00213
55106.856747	18.28734	0.00267
58043.863734	18.42718	0.00175
58068.809819	18.42294	0.00247
58089.731546	18.43190	0.00539
58089.843208	18.41577	0.00232
58116.705019	18.40079	0.00152
58117.703625	18.41215	0.00157
58129.814248	18.38507	0.00160
58156.798523	18.35187	0.00224

**Table B.3.** Radial velocity measurements for NGC4349 No. 127 obtained with HARPS and their instrumental error bars taken after JD-2 454 234. Older data is presented in Paper I. All data are relative to the solar system barycenter.

JD-2 400 000	<i>RV</i> [km s <sup>-1</sup> ]	Uncertainty [km s <sup>-1</sup> ]
54258.557903	-11.63628	0.00377
54293.529231	-11.67849	0.00298
54299.560888	-11.59373	0.00653
54319.473660	-11.63690	0.00664
54323.471152	-11.64945	0.00630
54342.475167	-11.64134	0.00325
54349.471386	-11.64002	0.00366
54481.827906	-11.47722	0.00356
54486.798601	-11.45261	0.00484
54523.803004	-11.31187	0.00511
54528.781327	-11.30763	0.00382
54553.727450	-11.30752	0.00384
54557.719105	-11.27895	0.00531
54567.681649	-11.25400	0.00460
54616.604433	-11.21941	0.00421
54620.561443	-11.21867	0.00733
54643.591230	-11.18338	0.00576
54878.815480	-11.48937	0.00312
54884.800643	-11.48642	0.00419
54915.765632	-11.59573	0.00384
54937.688513	-11.64931	0.00312
54939.702640	-11.67278	0.00507
54951.725335	-11.68053	0.00442
54993.573342	-11.73192	0.00349
55036.503475	-11.68980	0.00401